

Transverse Forward Physics at STAR (FMS)

April 29, 2010

Steve Heppelmann

Outline

- The Ideas
- Current STAR FPD/FMS Surprises
- Sensitivity of Future Measurements
- Photons
- Drell Yan
- Run 9 Setup Experience

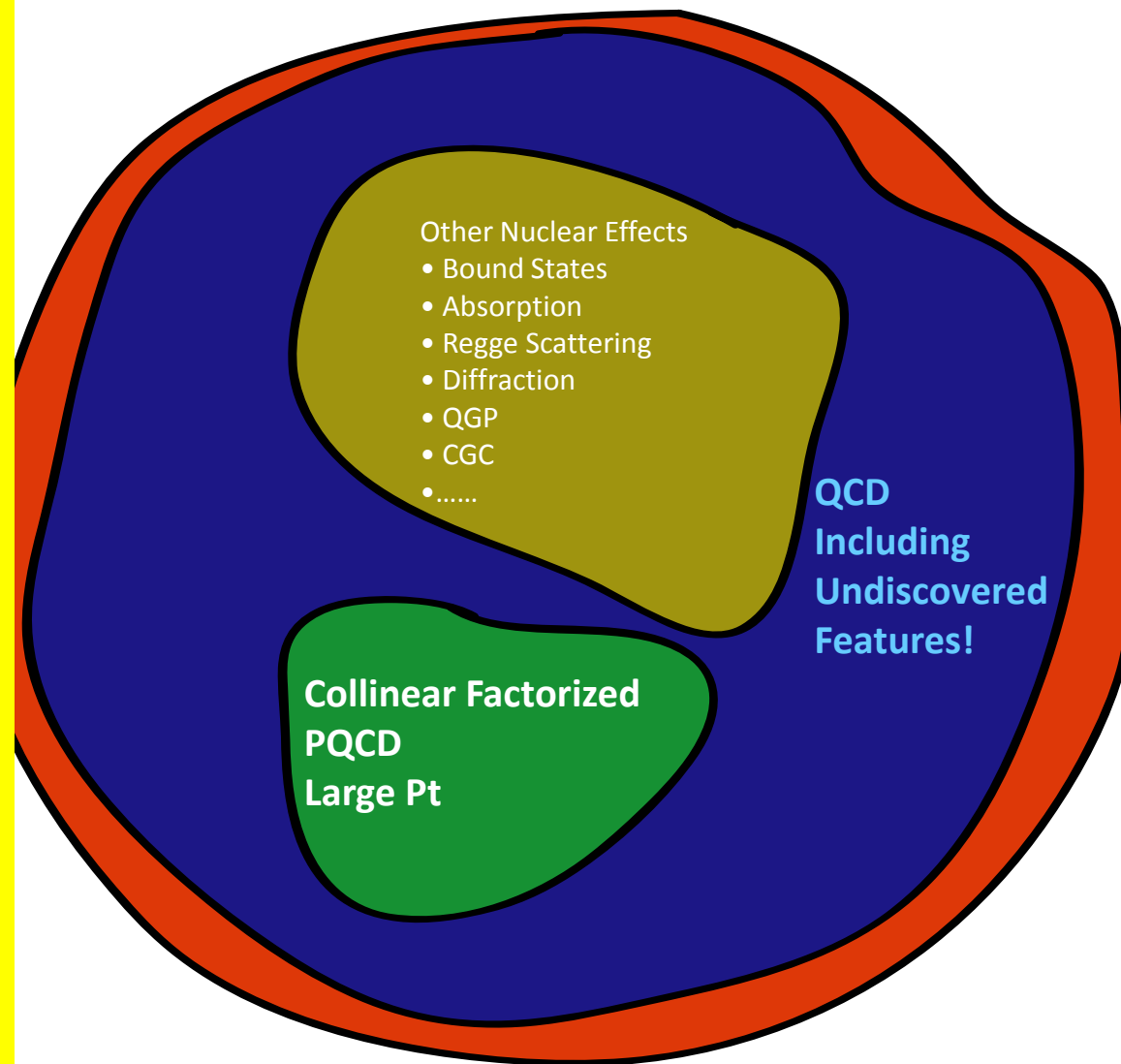
The Ideas: Transverse Forward Physics in STAR

- pQCD with collinear factorization should explain hard scattering with longitudinal polarization but transverse SSA vanishes in that framework.
 - Vanishing parton helicity flip
 - Leading order “real” scattering amplitudes.
 - Transverse single spin asymmetry (SSA) require helicity flip and imaginary amplitude.
 - SSA not allowed (in leading order/twist)
- **Expanding the scope of pQCD. Generalized extensions to collinear factorized.**
 - physics beyond collinear factorization
 - Generalized Parton Distributions
 - orbital angular momentum of partons
 - Sivers effect
 - Collins fragmentation
 - QCD Beyond Factorization

PQCD

Collinear Factorization

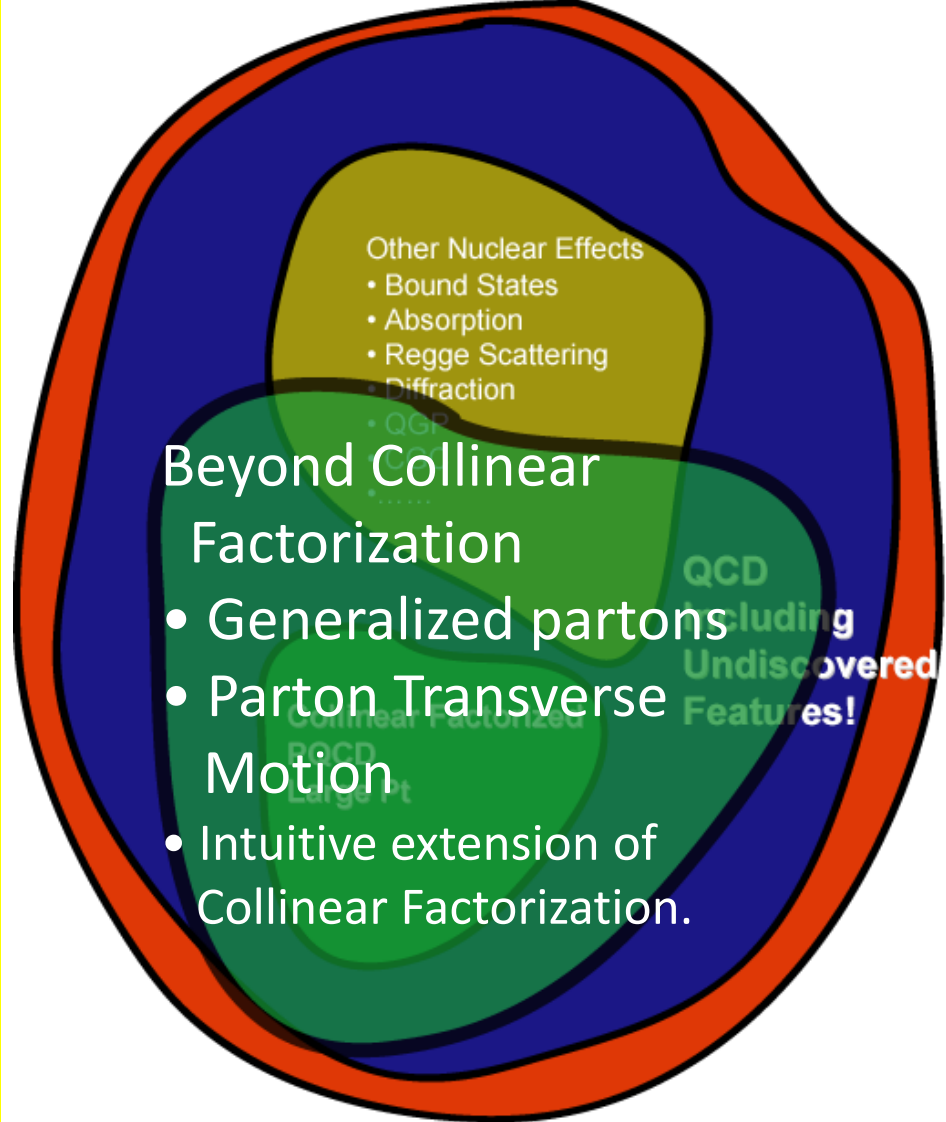
- **Gives meaning to quark and gluon**, the confined internal degrees of freedom (DOF) in QCD.
- Provides concrete connections between these **internal DOF** and **experimental observables**.
(Jets, some hadrons, photons)
- Gives an experimental connection to a **description of nucleon** and non-perturbative bound state (Nucleon parton densities) .
- Provides a recipe for **approximate calculation of cross sections** for certain interactions in certain kinematic regions.
- **Has a well defined kinematic region where calculations are most likely dependable.**



Strong Interactions

Generalized Factorization PQCD++

- Applies to a wider variety of experimental measurements.
- Gives similar meaning to quark and gluon, the confined internal degrees of freedom (DOF) in QCD. (same)
- Provides concrete connections between these internal DOF and experimental observables.
(Jets, some hadrons, photons) (same)
- Gives an experimental connection to a description of nucleon and non-perturbative bound state (Nucleon parton densities) . (same)
- Provides a recipe for approximate calculation of cross sections for certain interactions in certain kinematic regions??? (perhaps same)
- Has less clearly defined rules as to when calculations are most likely dependable.



Strong Interactions

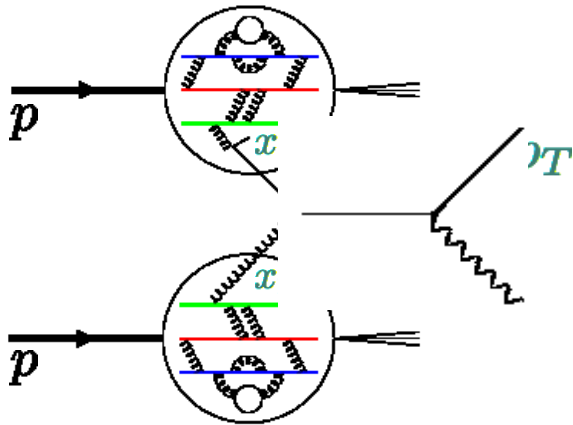
Collinear Factorization

Cross Section~ (Probability to select required parton A (x_1) from proton 1)
 x (Probability to select required parton B (x_2) from proton 2)

$$f_1(x_1) \underset{x_1 \rightarrow 1}{\sim} (1-x_1)^3$$

$$f_2(x_2) \underset{x_2 \rightarrow \text{small}}{\sim} 1$$

$$D_{parton}^{\pi^0} \underset{z \rightarrow 1}{\sim} (1-z)^1$$



$$d\sigma_{pp} \propto f_1 \otimes f_2 \otimes \sigma_h \otimes D_f^h$$

x (Probability that partons A+B => C + X)
 x (Probability that parton C Fragments into observed final state)

For Forward Production of Pi/Eta ..

$$\sigma(x) \propto \int_{x_f}^1 dz f_1 \left(x \sim \frac{x_F}{z} \right) \sigma_{parton} D_{parton}^{\pi^0}(z)$$

$$q(x) \sim (1-x)^3$$

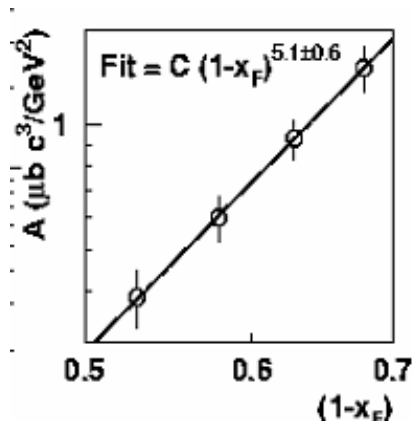
$$d(z) \sim (1-z)$$

$$\sigma(x) \propto (1-x_F)^5 + \text{Order}[(1-x_F)^6]$$

$$\sigma(x) \propto (1-x_F)^5$$

Forward Pi0 Cross Sections Scale Like seen in ISR.

At Large x_F (ie. $x_F > 0.4$), the Pi^0 fragment carries most of the of the jet momentum ($\langle z \rangle > 75\%$).



$$E \frac{d^3\sigma}{dp^3} \propto (1-x_F)^N p_T^{-B}$$

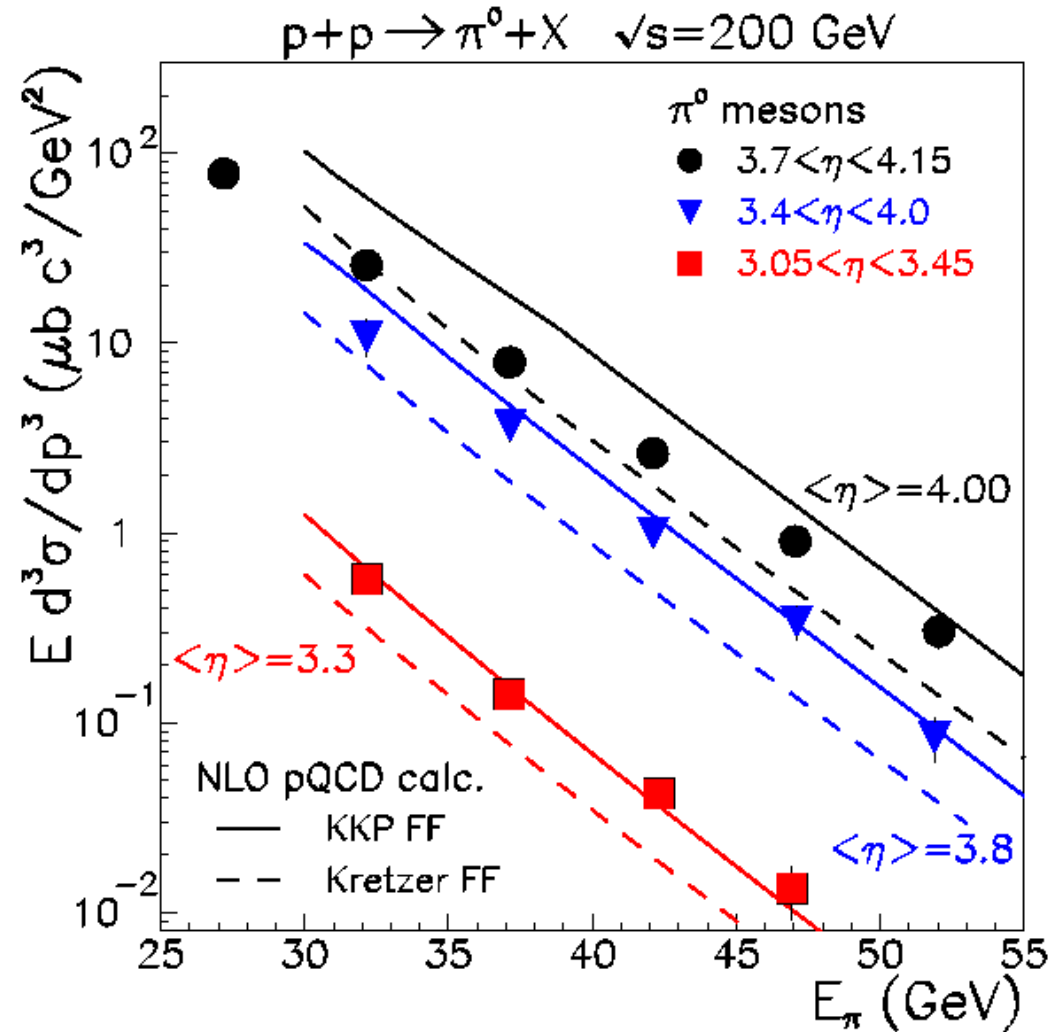
$$N \approx 5$$

$$B \approx 6$$

STAR Published
Result is similar to
to ISR analysis
J. Singh, et al Nucl.
Phys.
B140 (1978) 189.

$$\propto e^{-\frac{2(N+B)}{100}E} = e^{-\frac{2(11)}{100}E} = e^{-0.22E}$$

for $\{20 < E < 80\} \text{ GeV}$



Alternatives to Factorized PQCD Lead to Very Different Cross Sections

- Preliminary look at invariant cross section are likely consistent with conventional

$$\frac{(1 - x_F)^5}{p_T^6}$$

- In contrast, analysis of low p_T **Regge type processes** lead to to a different form for the dependence of the cross section on $(1-x_F)$ as Feynman x_F approach unity.

Regge Cross Section $\propto (1 - x_F)^2$

L.L.Frankfurt and M.I. Strikman, Vol. 94B2 Physics Letters, 28 July 1980.
and Private Communication.

Sivers Model

Difference Between pi0 and eta A_N ?

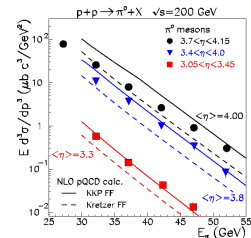
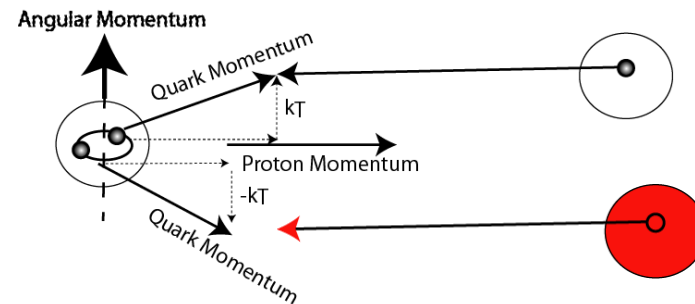
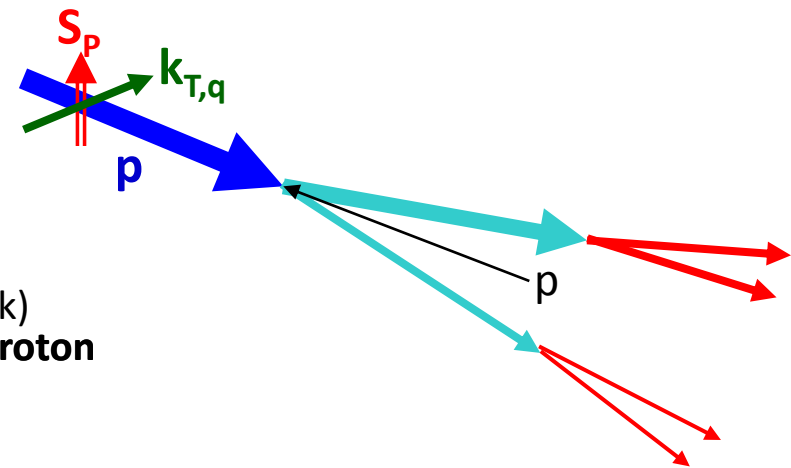
- A fast quark in the polarized proton (probably a u quark) has initial **transverse motion relative to the incident proton direction**. The sign of this transverse momentum is connected to the proton transverse spin.
- The jet has transverse momentum

$$P_T = p_T^{hard\ scattering} + k_T$$

- $\langle k_T \rangle$ changes sign if the spin and angular momentum is reversed.
- "T" symmetrical " $-k_T$ " amplitude absorbed as quark in one nucleon passes through gluon field of other nucleon. "Wilson Line"

Breaking of Factorization!!!!

- The jet fragments with large z to produce a **meson that is moving in the direction of the jet**, with nearly p_T of the jet.
- Dependence of **initial state** p_T upon proton spin leads to Sivers A_N .
- Shape of **cross section similar for pi0 and eta**.
- This situation should be the **same whether** the jet fragments into a **pi0 or an eta**.



$$A_N(P_T) = \frac{\sigma(p_T + \langle k_T \rangle) - \sigma(p_T - \langle k_T \rangle)}{2\sigma(P_T)}$$

$$\sim \frac{1}{\sigma} \frac{d\sigma}{dP_T} \langle k_T \rangle$$

Collins Model

k_T and thus A_N vanishes as Z approaches 1

- Consider large eta A_N (perhaps of order unity)
 $x_F \sim 0.75$, $Z \sim .9$ and $p_T \sim 3.9$ GeV/c.
- Any associated jet fragments will carry limited transverse momentum,

$$k_T \sim \frac{(1-Z)p_T}{2}$$

- If the cross section is given by $\frac{(1-x_F)^5}{p_T^6}$

- The Maximal asymmetry from fragmentation $p_T \rightarrow p_T + \sin(\phi)k_T$

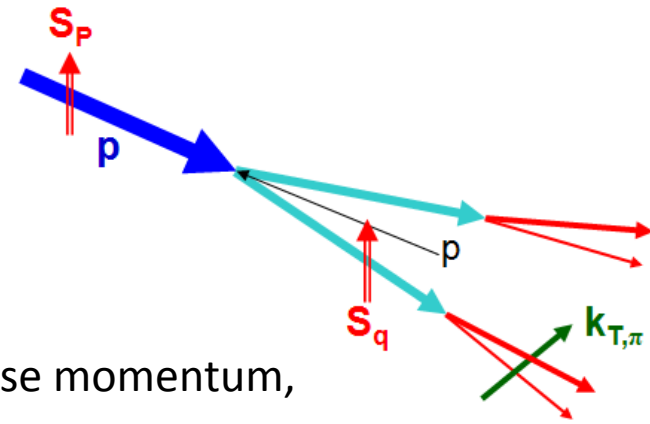
$\phi =$ fragmentation azimuthal angle from spin direction

- Leads to an **extreme limit** for A_N from fragmentation,

$$A_N < \frac{6k_T}{p_T} \sim 3(1-Z) \sim .6$$

This is the most extreme case including

- 100% transverse parton polarization
- the maximum possible Collins Fragmentation function.



Current STAR+FPD/FMS Surprises

P_t Dependence in Calculations of A_N

•Sivers Effect / Collins Effect

•introduce transverse spin dependent offsets in transverse momentum

•independent of the hard scattering (definition of factorization).

$$P_T \Rightarrow P_T \pm k_T$$

“ \pm ” depending on the sign of proton transverse spin direction. Using our (STAR) measured cross section form:

$$d\sigma^\uparrow \propto \frac{1}{(P_T - k_T)^6} \quad d\sigma^\downarrow \propto \frac{1}{(P_T + k_T)^6}$$

$$A_n \equiv \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow} = \frac{6k_T}{P_T} + O\left(\frac{k_T}{P_T}\right)^2$$

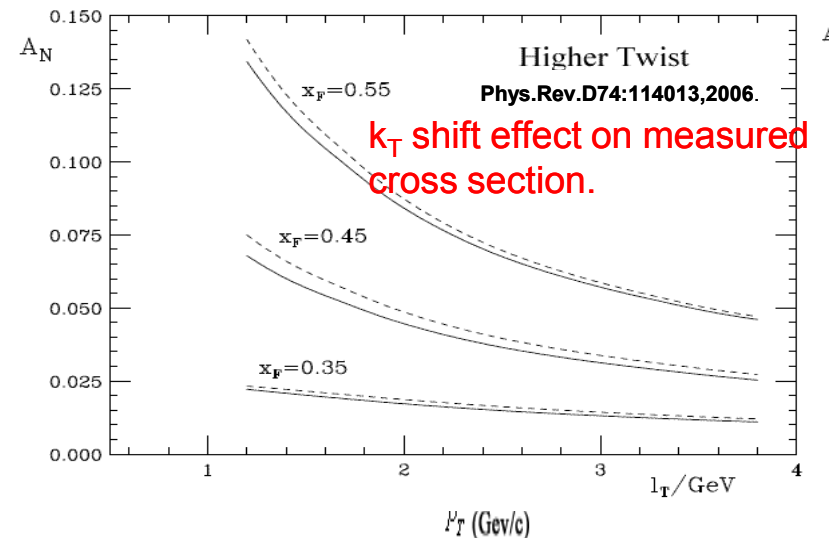
Higher Twist Effects:

Qiu and Sterman

Kouvaris et. al. **Phys.Rev.D74:114013,2006.**

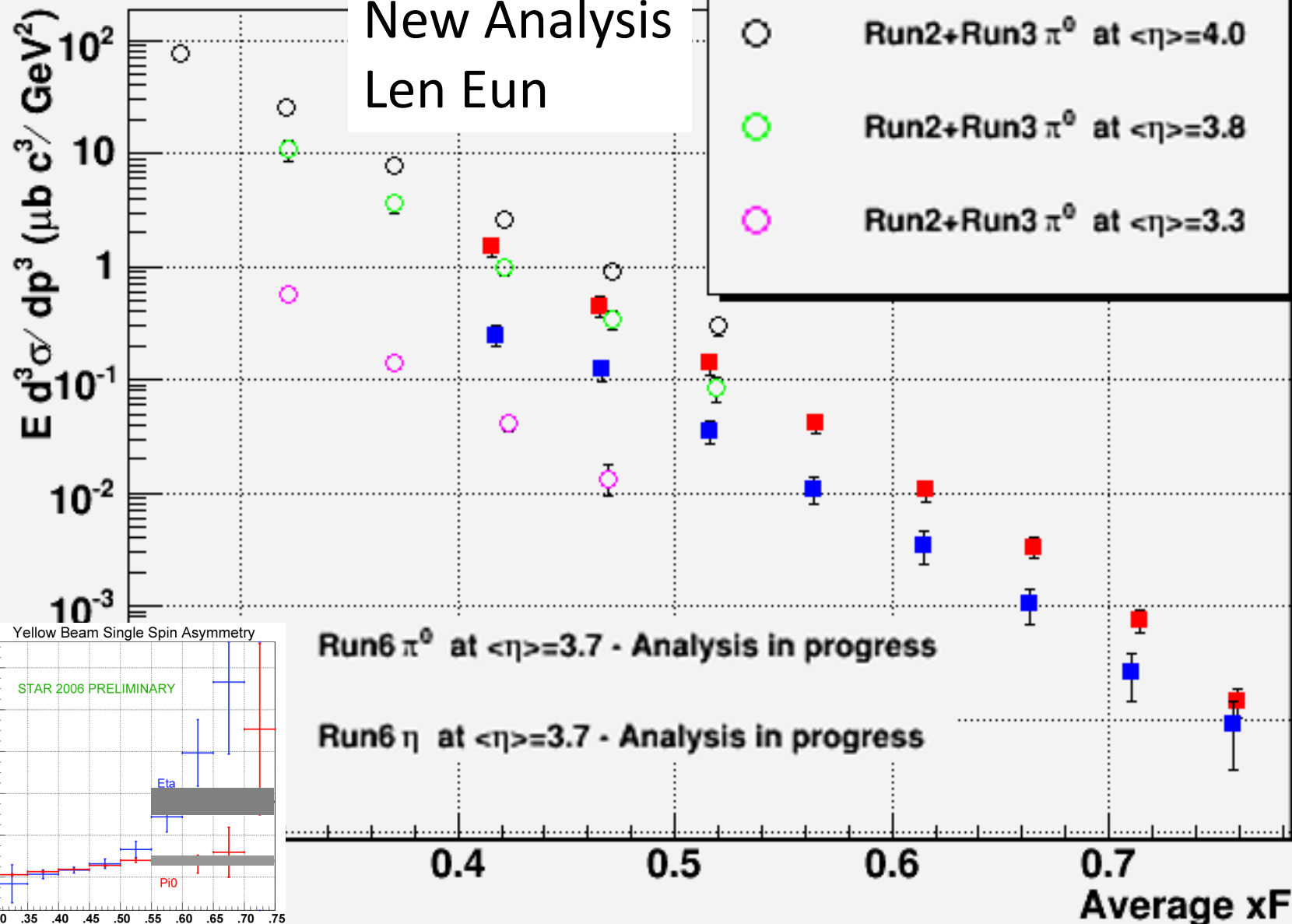
A_N Fall as $1/P_T$ as required by definition of higher twist.

All of these models lead to
 $A_N \sim \propto 1/P_T$

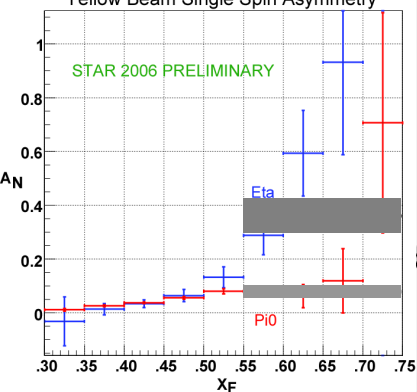


Preliminary Invariant Cross-sections - Calibration Not Finalized

New Analysis
Len Eun



Yellow Beam Single Spin Asymmetry



Comparison between η production and π^0 production?

- **Gluons** or η has **Isospin $I=0$** .
- **u quark** has **Isospin $I=1/2$**
- π^0 has **Isospin $I=1$** .
- But we expect both mesons to come from fragmentation of quark jets.

$$I = 0 \quad \left\{ \begin{array}{l} \eta \simeq \frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d} - s\bar{s}) \\ \eta' \simeq \frac{1}{\sqrt{6}}(u\bar{u} + d\bar{d} + 2s\bar{s}) \end{array} \right.$$

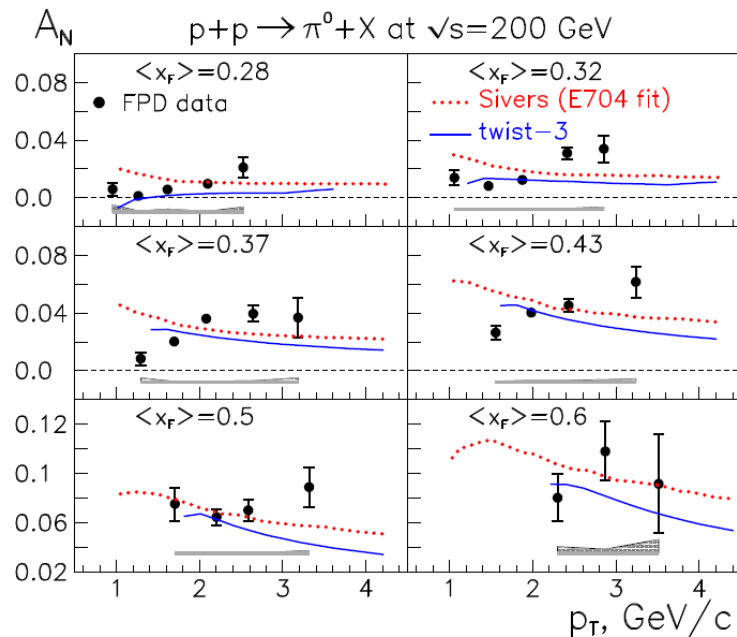
$$I = 1 \quad \left\{ \begin{array}{l} \pi^0 = \frac{1}{\sqrt{2}}(u\bar{u} - d\bar{d}) \end{array} \right.$$

*Assume η, η' mixing angle: $\theta_p \sim -19.5^\circ$

- **For Sivers Effect:** Asymmetry is in the jet and should not depend on the details of fragmentation.
- **For Collins Effect:** Asymmetry reflects fragmentation of the quark jet into a leading η or π^0 meson. Differences in fragmentation could relate to:
 - Mass differences?
 - Isospin differences?
 - Role of Strangeness?
 - **But Collins Effect Should be suppressed when $Z \rightarrow 1$**

For Fixed X_F , the asymmetry A_N does not fall with P_T as predicted by models.

- NLO PQCD **does describe** the size and shape of this forward pp cross section.
- Model calculations (Sivers, Collins or twist-3) **can explain** the X_F dependence of A_N .
- Flat or increasing dependence of A_N on P_T**



U. D'Alesio, F. Murgia, Phys. Rev. D **70**, 074009 (2004).

J. Qiu, G. Sterman, Phys. Rev. D **59**, 014004 (1998).

Theory Score Card For Factorized QCD Picture for Pi & Eta Transverse A_N

✓ Cross Section for Pi0 agrees with PQCD (Normalization and Shape)

✓ Dependence of cross section on X_F and P_T may be similar for Pi0 and Eta at large X_F as expected.

✓? Ratio Eta/Pi0 nominal 40% - 50%
Yet to be determined.

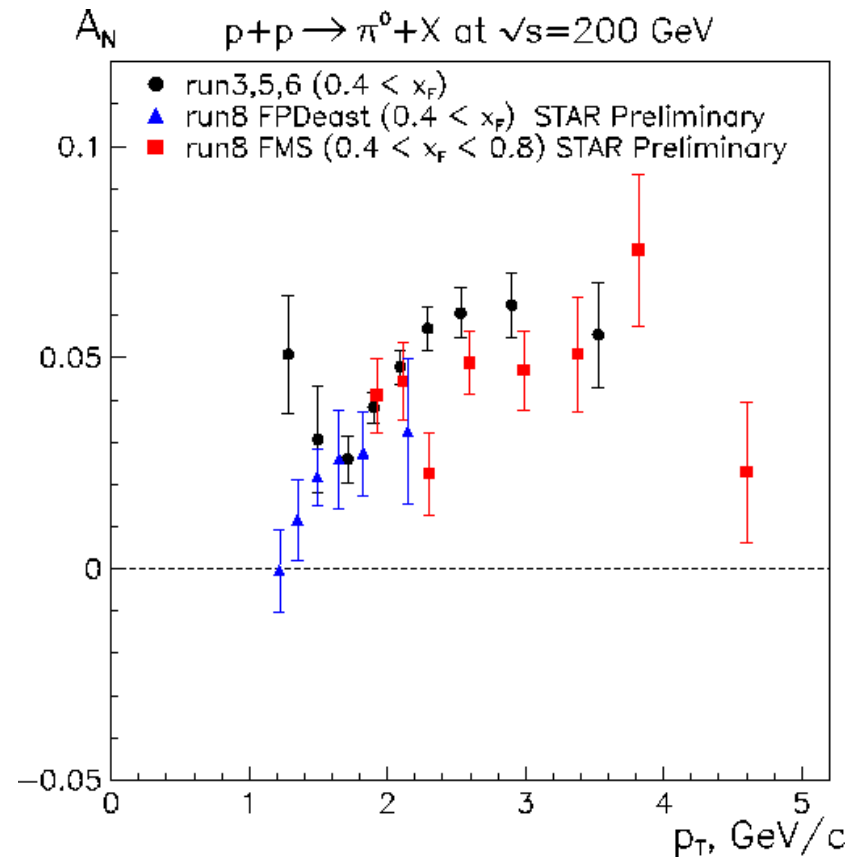
✗ Pt Dependence of Pi0 A_N .

Inconsistent with $A_N \sim 1/p_T$.

Can a large difference in asymmetry between Pi0's and Eta's be understood in either Collins or Sivers Model?

With FMS, STAR has Expanded Rapidity Coverage $-1 < Y < 4.2$

STAR Forward Meson Spectrometer
 $2.5 < Y < 4.0$



**arXiv:0901.2763 +
A.Ogawa @CIPANP09**

Sensitivity of Future Measurements

200 GeV Transverse Spin Program

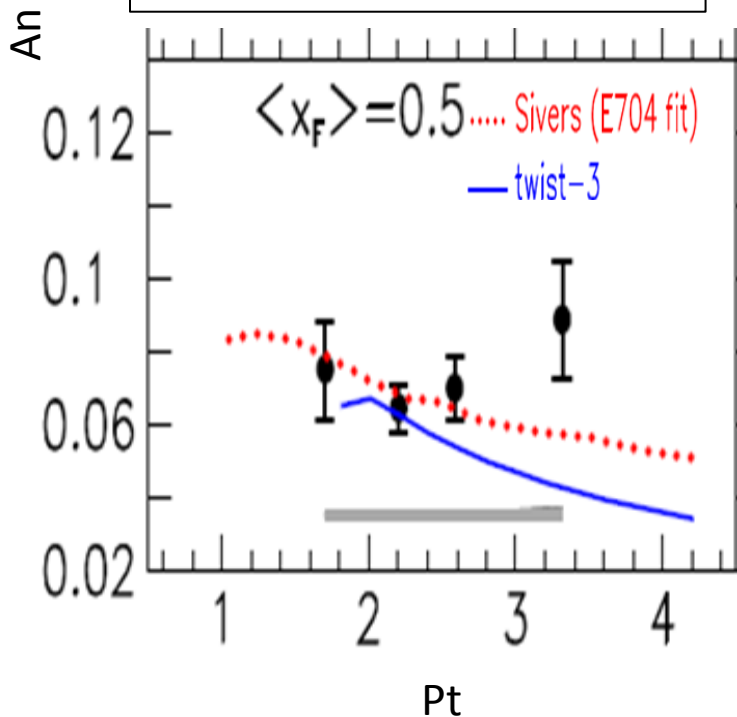
\sqrt{s}	Transverse Process	Collins	Sivers	Sivers SIDIS Sign Change	$\int L$	Detectors
200 GeV	$p^\uparrow + p \rightarrow \pi^0 + X$	*	*		30 pb^{-1}	FMS
	$p^\uparrow + p \rightarrow \eta + X$	*	*			
	$p^\uparrow + p \rightarrow \gamma + X$		*	*		
	$p^\uparrow + p \rightarrow \pi^0 + \pi^0 + X$	*	*			
	$p^\uparrow + p \rightarrow \gamma + \pi^0 + X$		*	*		
	$p^\uparrow + p \rightarrow \pi^0 + \pi^0 + X$	*	*			FMS+EMC
	$p^\uparrow + p \rightarrow jet + X$	*	*			FMS+HCAL
	$p^\uparrow + p \rightarrow jet + \pi^0 + X$	*	*			
	$p^\uparrow + p \rightarrow \Lambda + X$	*	*			
	$p^\uparrow + p \rightarrow jet + jet + X$	*	*			FMS+EMC
	$p^\uparrow + p \rightarrow \gamma + jet + X$		*	*		HCAL +Tracking

500 GeV Transverse Spin Program

\sqrt{s}	Transverse Process	Collins	Sivers	Sivers SIDIS Sign Change	$\int L$	Detectors
500 GeV	$p^\uparrow + p \rightarrow \pi^0 + X$	*	*		20 pb^{-1}	West FMS East FPD +Shower Max
	$p^\uparrow + p \rightarrow \eta + X$	*	*			
	$p^\uparrow + p \rightarrow \gamma + X$		*	*		
	$p^\uparrow + p \rightarrow \pi^0 + \pi^0 + X$	*	*			
	$p^\uparrow + p \rightarrow \gamma + \pi^0 + X$		*	*		
	$p^\uparrow + p \rightarrow \pi^0 + \pi^0 + X$	*	*			FMS+EMC
	$p^\uparrow + p \rightarrow jet + X$	*	*			FMS+HCAL
	$p^\uparrow + p \rightarrow jet + \pi^0 + X$	*	*			
	$p^\uparrow + p \rightarrow \Lambda + X$	*	*			
	$p^\uparrow + p \rightarrow jet + jet + X$	*	*			FMS+EMC HCAL +Tracking
	$p^\uparrow + p \rightarrow \gamma + jet + X$		*	*		
	$p^\uparrow + p \rightarrow e^+ + e^- + X$		*	*	250 pb^{-1}	FMS+EMC +Tracking +PID

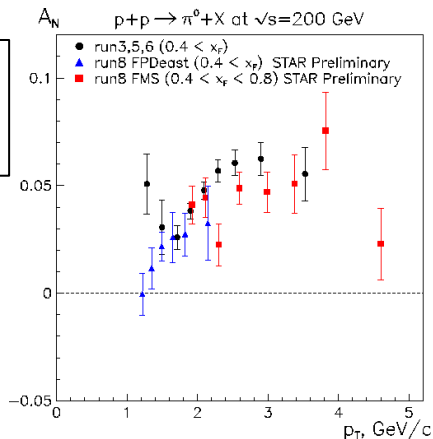
$$A_N \pi^0$$

Run 6 FPD Pt Dependence
 $X_F \sim 0.5$

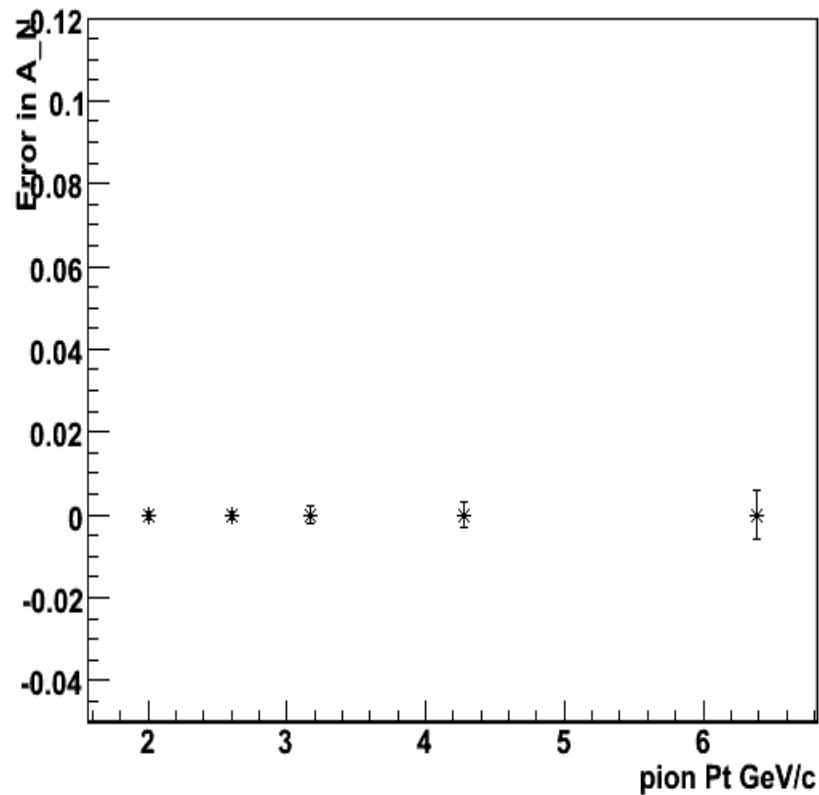


Run 6 FPD+ **Run 8 FMS**
Pt Dependence $X_F > 0.4$

Errors for Projected
FMS Pt dependence
 $0.5 < X_F < 0.55$

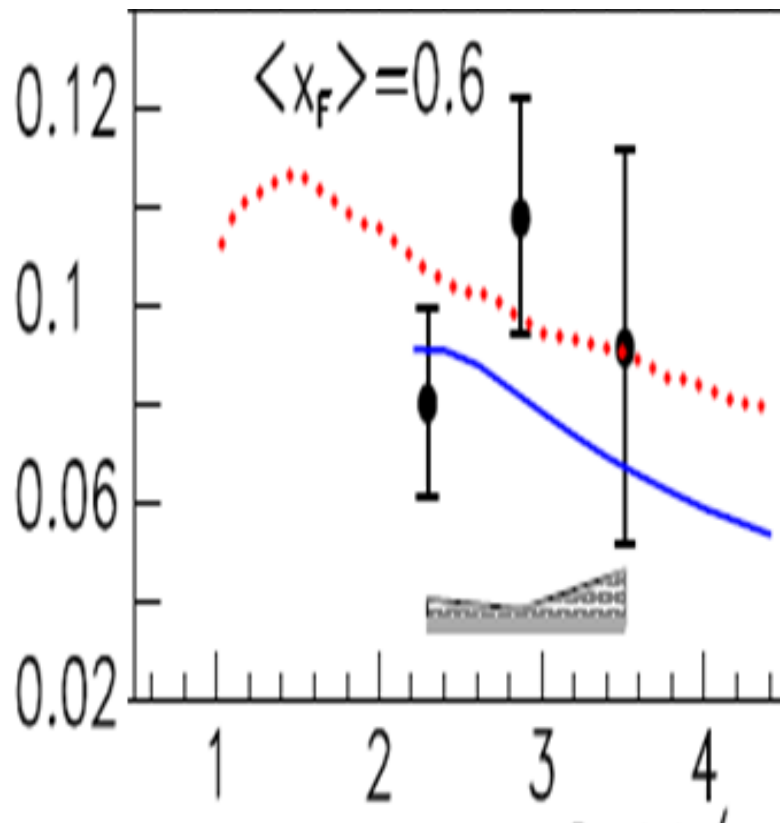


Projected Error in $\pi^0 A_N$: 30 pb⁻¹(-1) root_s=200 .50<X_F<.55



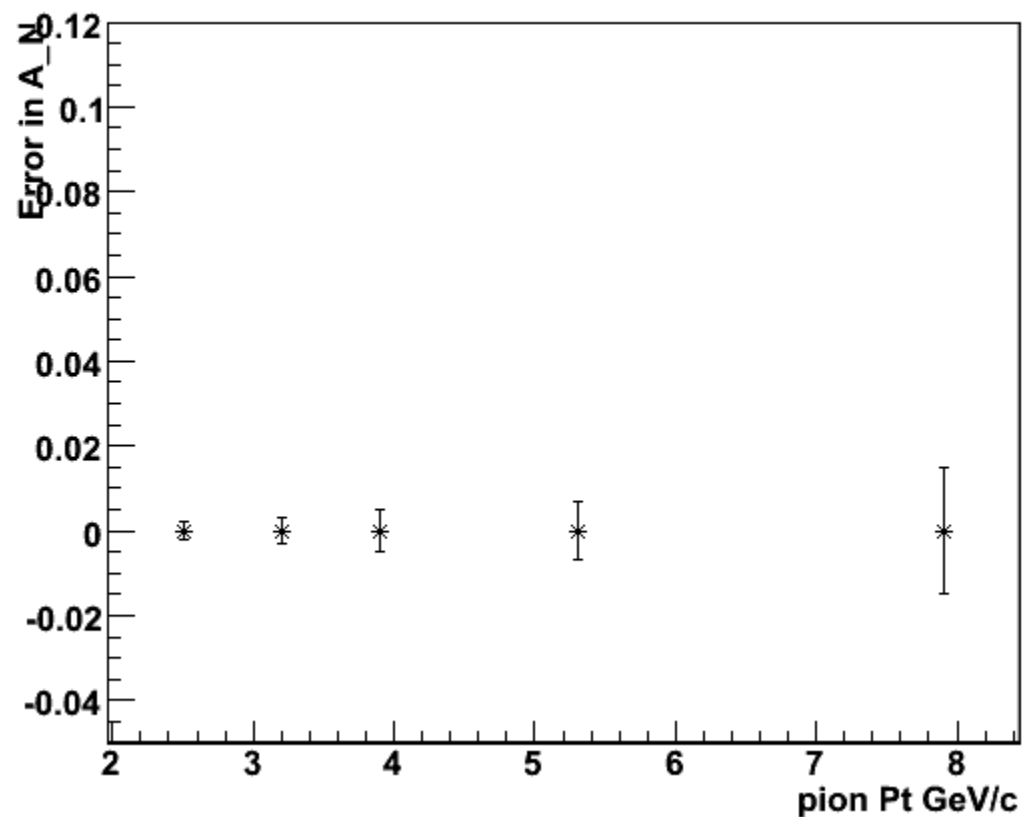
$$A_N \pi^0$$

Run 6 FPD Pt Dependence
 $X_F \sim 0.6$



Errors for Projected
 FMS Pt dependence
 $0.6 < X_F < 0.65$

Projected Error in $\pi^0 A_N$: 30 pb⁻¹ root_s=200 .6<X_F<.7





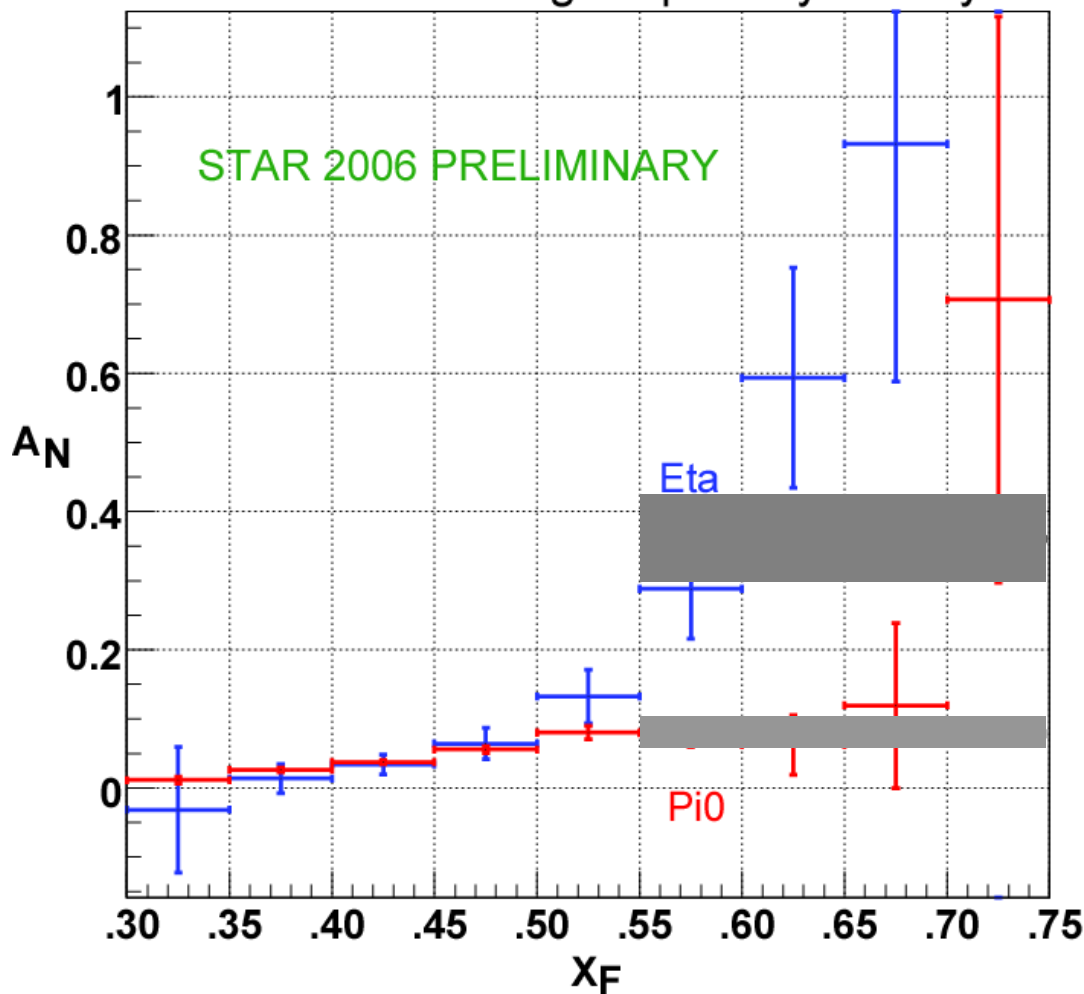
$A_N(x_F)$ in π^0 and Eta Mass Regions

$$p^\uparrow + p \rightarrow M + X$$

$$M \rightarrow \gamma + \gamma \quad \sqrt{s} = 200 \text{ GeV}$$

Yellow Beam Single Spin Asymmetry

STAR 2006 PRELIMINARY



1. $N_{\text{photon}} = 2$
2. Center Cut (η and ϕ)
3. **Pi0** or **Eta** mass cuts
4. Average Yellow Beam Polarization = 56%

$$.55 < X_F < .75$$

$$\langle A_N \rangle_\eta = 0.361 \pm 0.064$$

$$\langle A_N \rangle_\pi = 0.078 \pm 0.018$$

For $.55 < X_F < .75$ the asymmetry in the η mass region is greater than 5 sigma above zero, and about 4 sigma above the asymmetry in the π^0 mass region.

Yellow Beam Single Spin Asymmetry

Run 6 FPD Result

$$.55 < X_F < .75$$

$$\langle A_N \rangle_\eta = 0.361 \pm 0.064$$

$$\langle A_N \rangle_\pi = 0.078 \pm 0.018$$

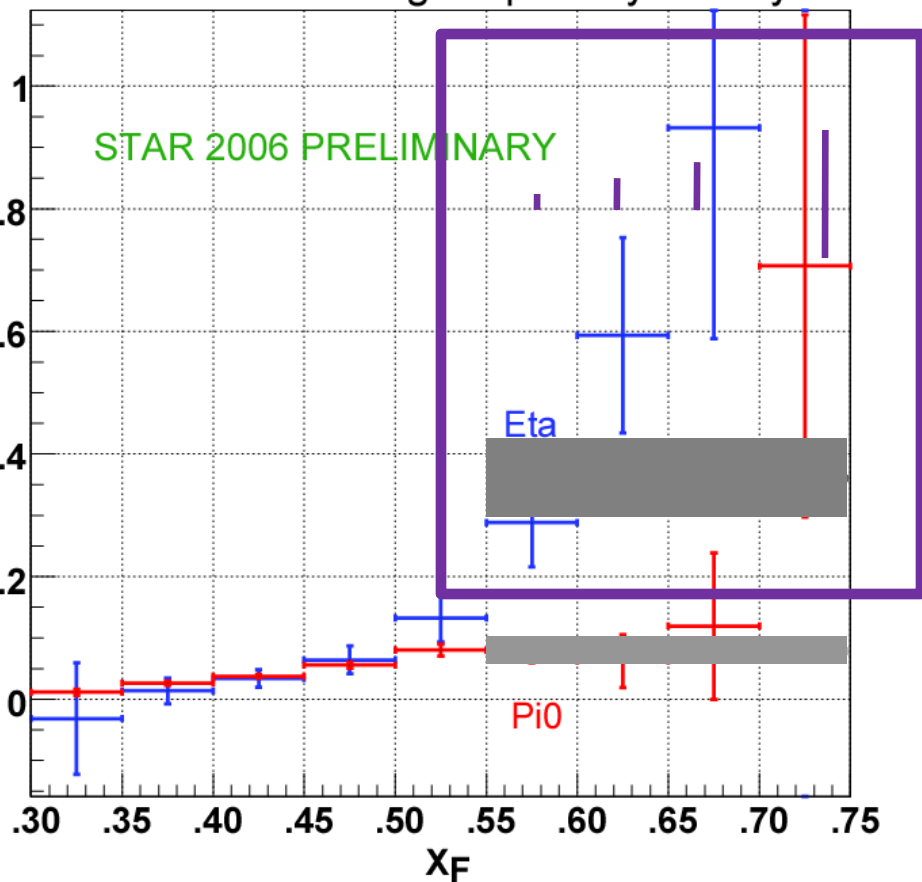
Projected Errors

For Eta A_N

200 GeV

30 pb⁻¹

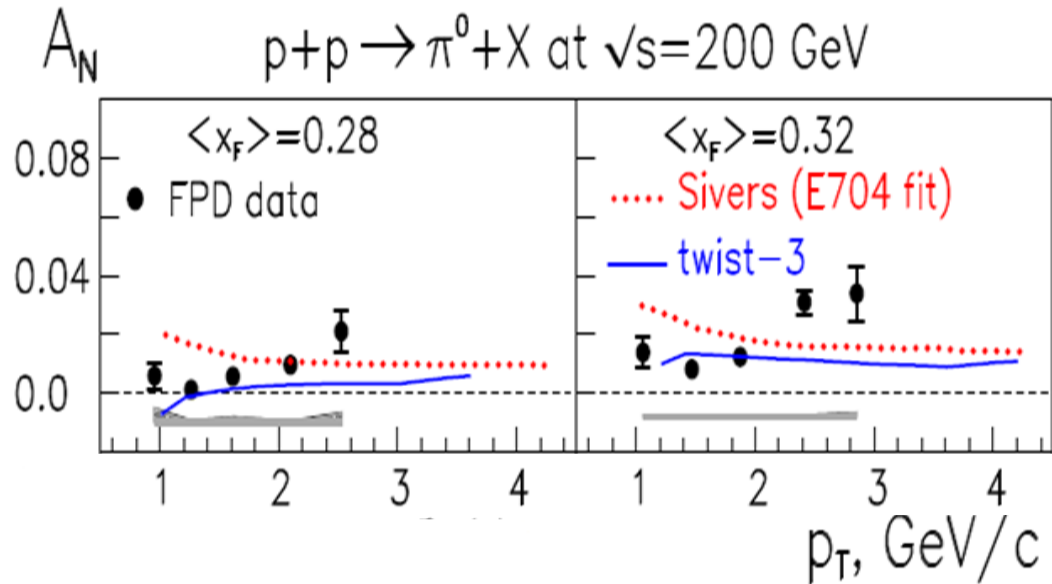
A_N



	$.50 < x_F < .55$	$.55 < x_F < .60$	$.60 < x_F < .70$	$.70 < x_F < .90$
$2.6 < Y < 3.0$	0.043	0.070	0.105	0.333
$3.0 < Y < 3.4$	0.019	0.031	0.047	0.149
$3.4 < Y < 3.6$	0.015	0.024	0.037	0.116
$3.6 < Y < 3.8$	0.010	0.016	0.025	0.078
$3.8 < Y < 4.1$	0.005	0.008	0.012	0.038

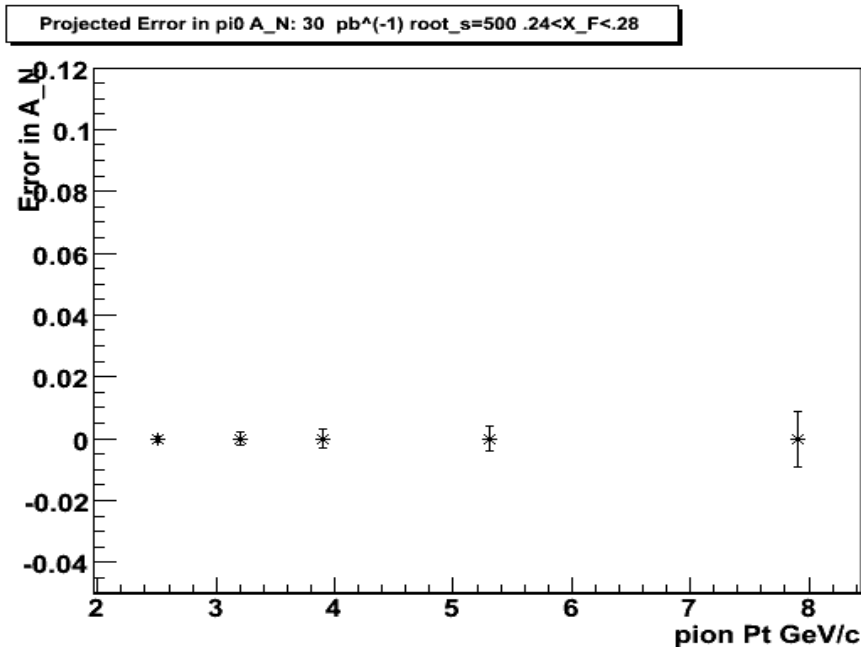
Table 5: Estimated error in A_N for η mesons for 200 GeV 30pb⁻¹ Data Set

Comparison between 200 GeV Measurement and 500 GeV Projections



Run 6 FPD Published
 A_N measurement

$\sqrt{s} = 200$ GeV



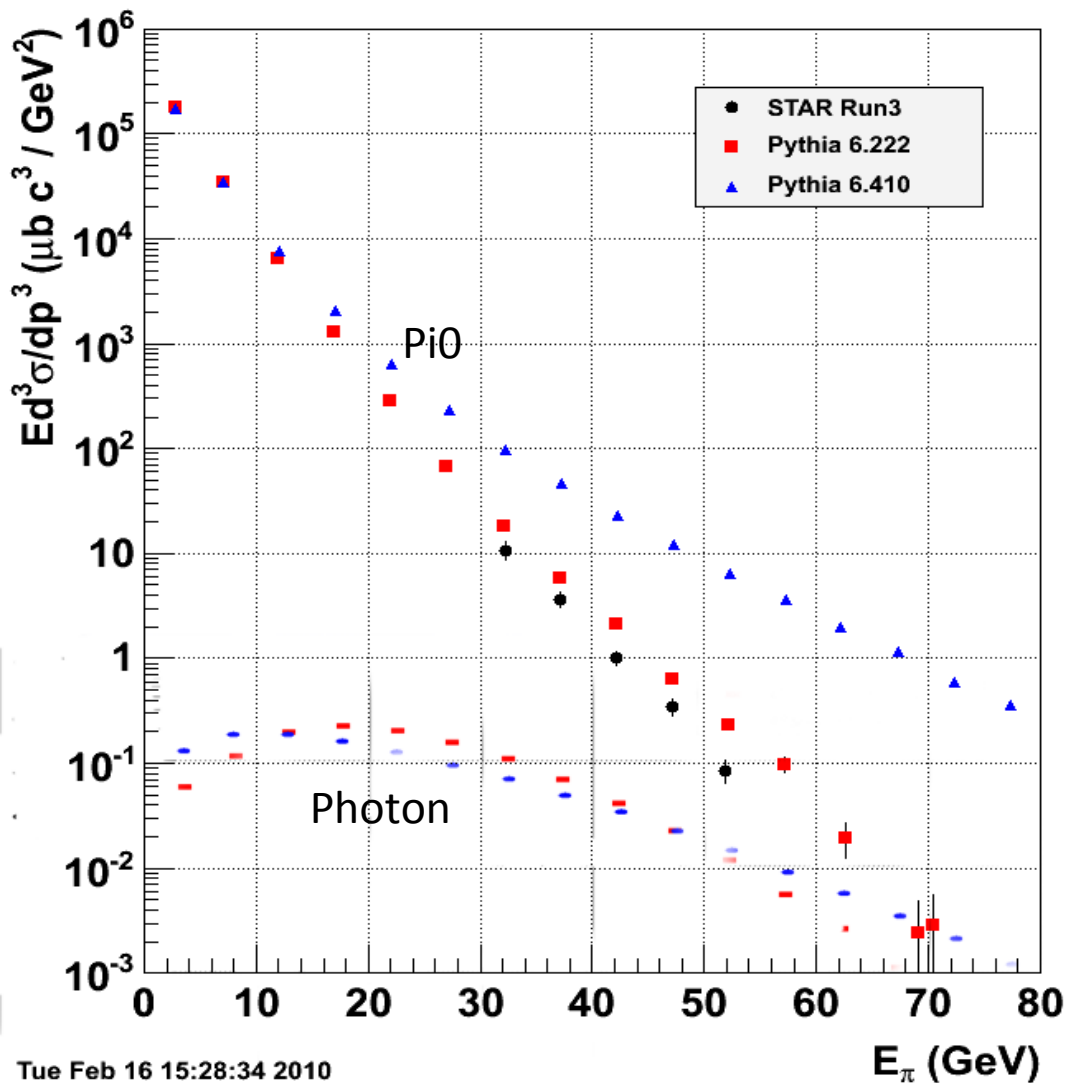
Projections for A_N
statistical errors

$\sqrt{s} = 500$ GeV
 20 pb^{-1}

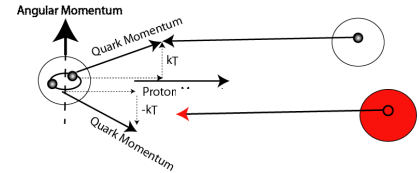
Photons

What Pythia says
For π^0 and γ

STAR data



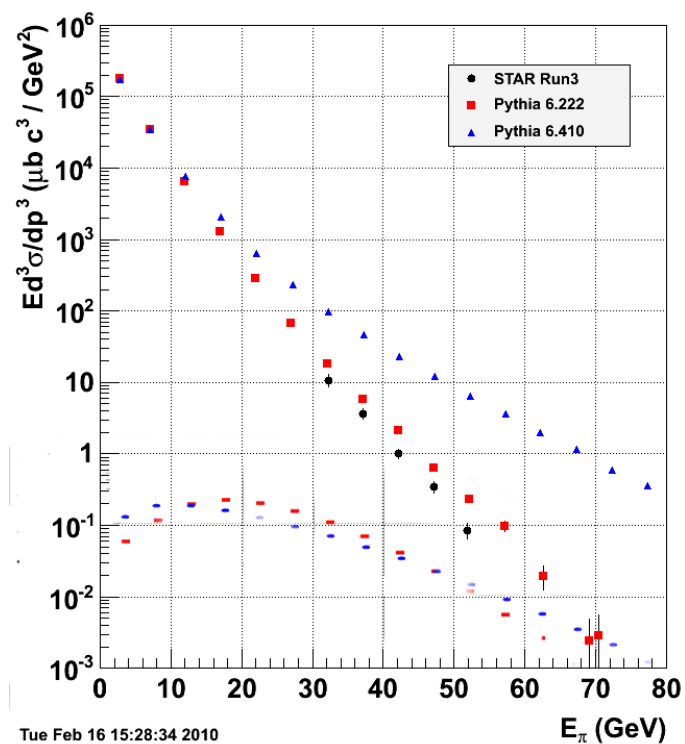
Direct Photon A_N Measurement



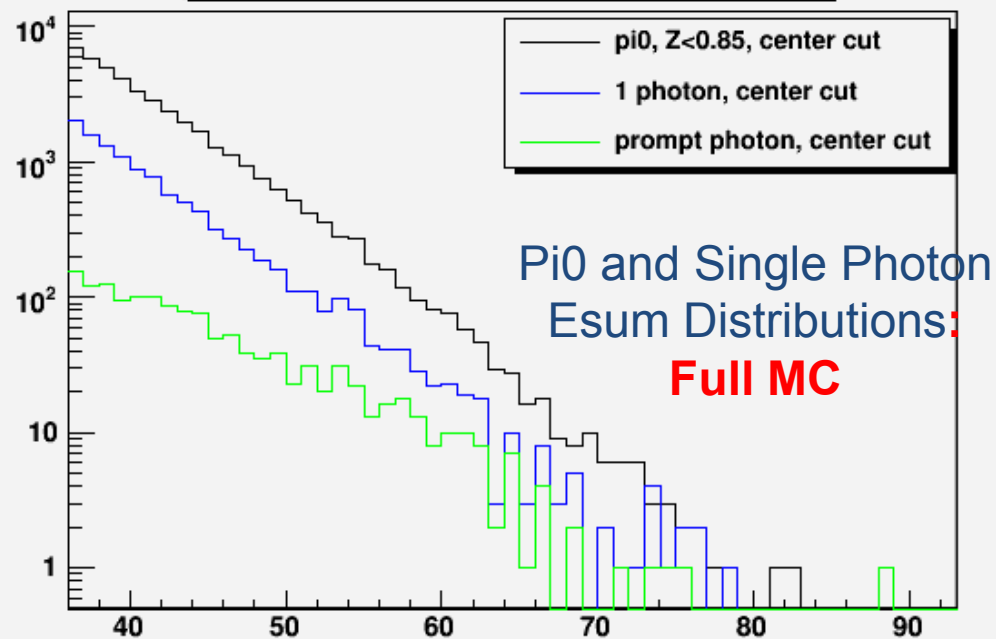
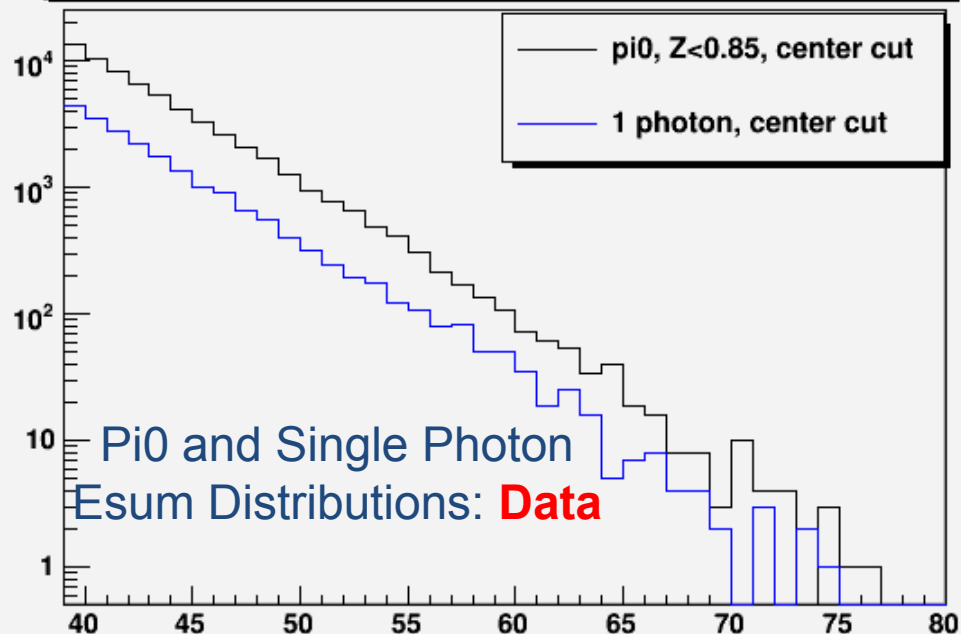
- Predicted violation of factorization
 - If Sivers is mechanism: **a sign change** is predicted between Direct Photon and DIS.
 - No Collins effect in Direct Photon A_N .
- **Measurement of predicted sign change vs A_N in DIS is a milestone goal from Nuclear Science Advisory Committee.**
- For $X_F > .5$, single photon cross section similar to π^0 cross section (see previous error estimates).
- Separation of 1 photon from 2 photon clusters based upon shower shape.
- Statistical errors similar to that for π^0 .
- Full errors dominated by background subtraction. (π^0 and η).

FPD Run 6 DATA and Simulations

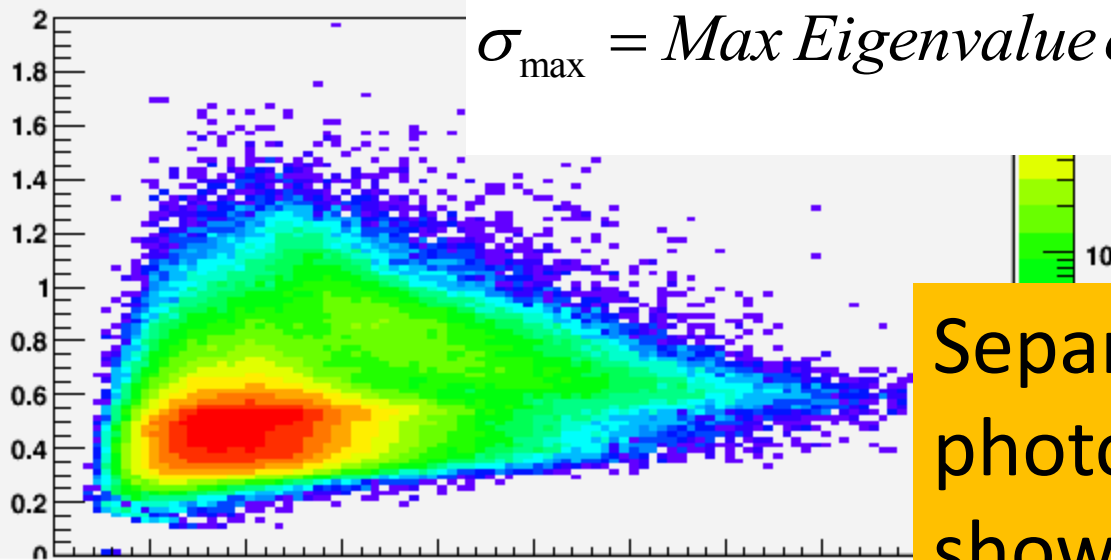
2 Photon pi0 events and 1 photon events from Len Eun



Run6T data: Summed Energy Distributions for pi0 and 1 photon events



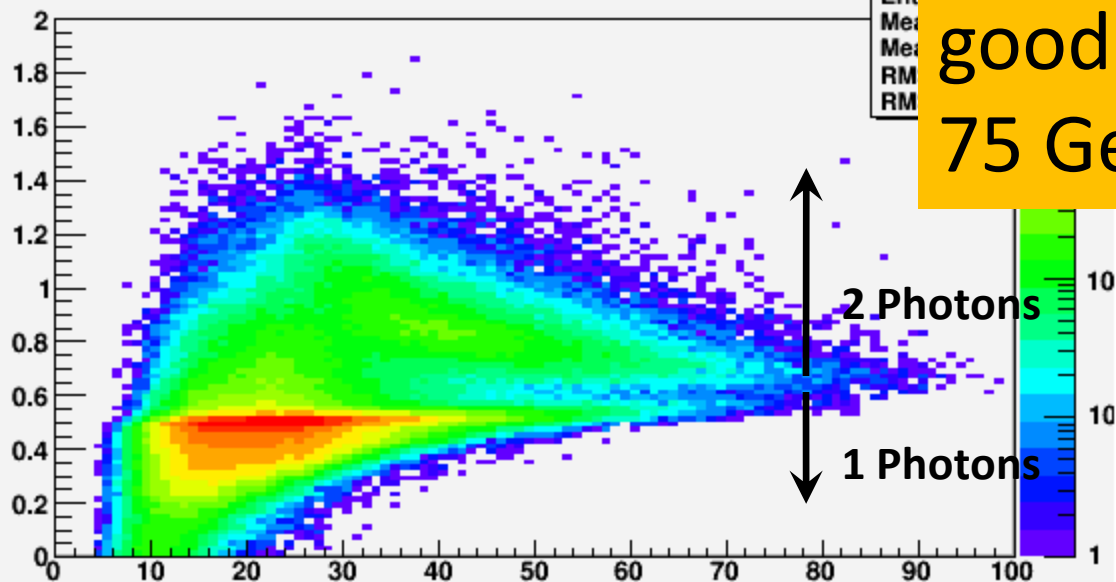
Esum vs. SigmaMax, Inner, Old



$$\sigma_{\max} = \text{Max Eigenvalue of } \begin{bmatrix} \Delta\sigma_x^2 & \Delta\sigma_x\Delta\sigma_y \\ \Delta\sigma_y\Delta\sigma_x & \Delta\sigma_y^2 \end{bmatrix}$$

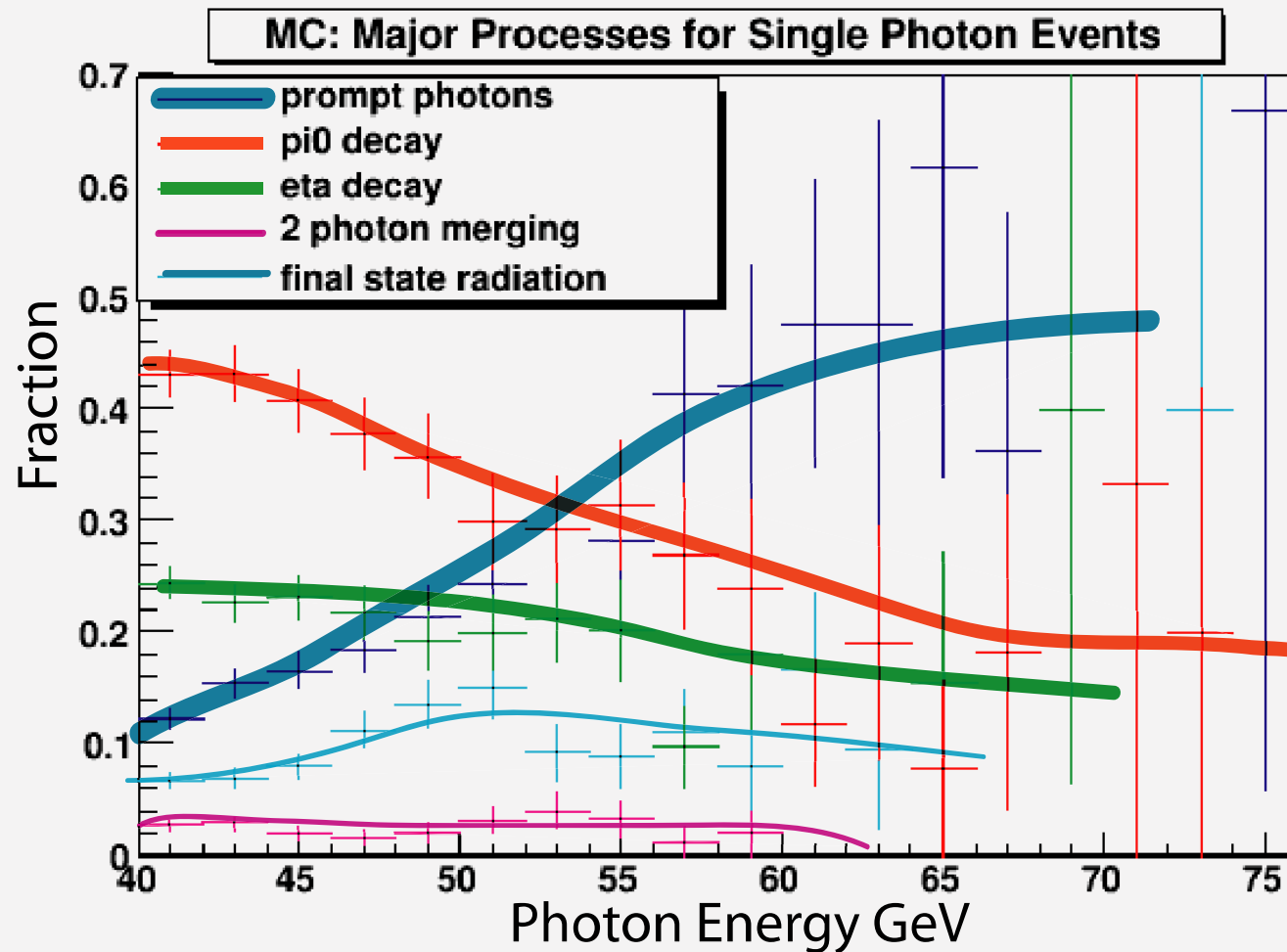
Separation of 1 vs 2 photons based on shower shape good to beyond 75 GeV

Esum vs. SigmaMax, Inner, New



Separation of single photon from two photon cluster based upon shower shape.

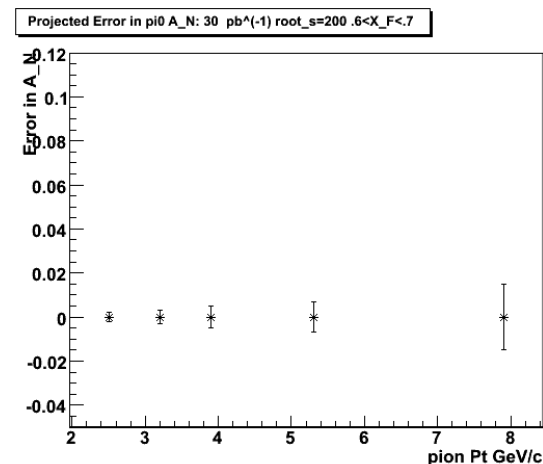
Energy GeV



Extracting Photon A_N

- FMS Run 9 data for energy > 65 GeV is approximately consistent with Pythia 6.222. This FMS data has little overlap with published FPD measurement.

- 30 pb of 200 GeV should produce
 - 50K π^0 with $E > 65$ GeV ; $3.6 < Y < 3.7$
 - 20K 1 photon events for a 1% measurement of A_N
 - including 50% real direct photons
 - including 25% photons from π^0
 - including 25% photons from η



- Determination of Single Photon Asymmetry Must be associated with a **comparable determination of the Eta and Pi0 asymmetries** at high energy.

$$\Delta A = \frac{N^\gamma A^\gamma + N^{\pi^0} A^{\pi^0} + N^\eta A^\eta}{N^\gamma + N^{\pi^0} + N^\eta} = 0.5 A^\gamma + 0.25 A^{\pi^0} + 0.25 A^\eta$$

	$.50 < x_F < .55$	$.55 < x_F < .60$	$.60 < x_F < .70$	$.70 < x_F < .90$
$2.6 < Y < 3.0$	0.006	0.010	0.015	0.047
$3.0 < Y < 3.4$	0.003	0.004	0.007	0.021
$3.4 < Y < 3.6$	0.002	0.003	0.005	0.016
$3.6 < Y < 3.8$	0.001	0.002	0.003	0.011
$3.8 < Y < 4.1$	0.001	0.001	0.002	0.005

Table 3: Projections for π^0 A_N : Entrys indicate the expected error in the transverse SSA (ΔA_N) for integrated luminosity of 30 pb^{-1} at $\sqrt{s} = 200 \text{ GeV}$. The values in table columns are for indicated vales of Feynman x_F and the rows are for indicated values of pseudorapidity.

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$3.8 < Y < 4.1$	0.005	0.008	0.012	0.038

Table 5: Estimated error in A_N for η mesons for 200 GeV 30 pb^{-1} Data Set

	$.2 < x_F < .22$	$.22 < x_F < .24$	$.24 < x_F < .28$	$.28 < x_F < .36$	$.36 < x_F < .4$
$2.6 < Y < 3.0$	0.004	0.006	0.009	0.030	0.188
$3.0 < Y < 3.4$	0.002	0.003	0.004	0.013	0.084
$3.4 < Y < 3.6$	0.001	0.002	0.003	0.010	0.065
$3.6 < Y < 3.8$	0.001	0.001	0.002	0.007	0.044
$3.8 < Y < 4.1$	0.0004	0.001	0.001	0.003	0.022

Table 7: Indicated the estimated error in the transverse SSA (A_N) for production of π mesons with integrated luminosity of 20 pb^{-1} at $\sqrt{s} = 500 \text{ GeV}$.

Drell Yan

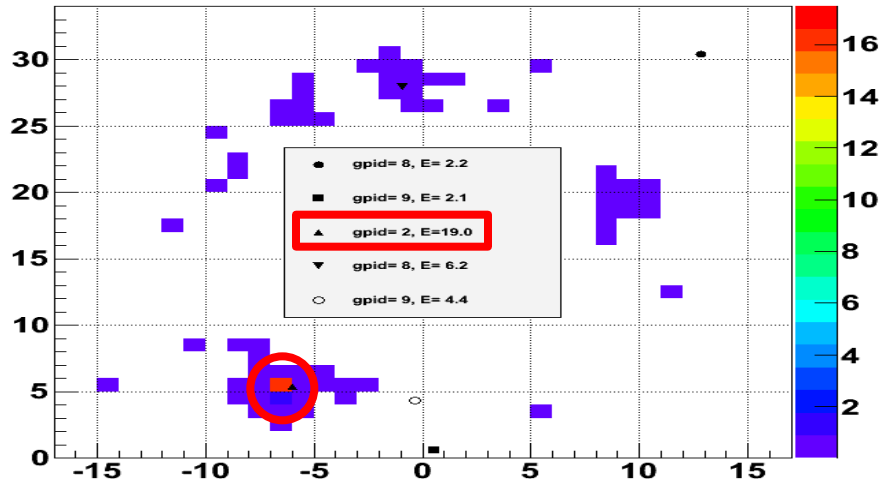
Drell Yan 500 GeV

- FMS $E_1 > 20$ GeV & $E_2 > 20$ GeV & $Mass > 4$ GeV;
- 250 pb⁻¹ -> 60k DY pairs -> $\Delta A_N \sim .01$ (statistical only)
- If neutral particles are rejected, the hadronic background due to hadronic energy deposited in the FMS may be comparable to the DY signal.
- Tests of backgrounds probably required!
 - Further background suppression is possible if magnetic charge sign and magnetic momentum determination is available. **30% momentum measurement on each of two tracks -> $\sim 10^{-3}$ suppression of background.**
 - Further background suppression is possible if charge transition radiation detector is available. **$\sim 10^{-3}$ for two tracks.**

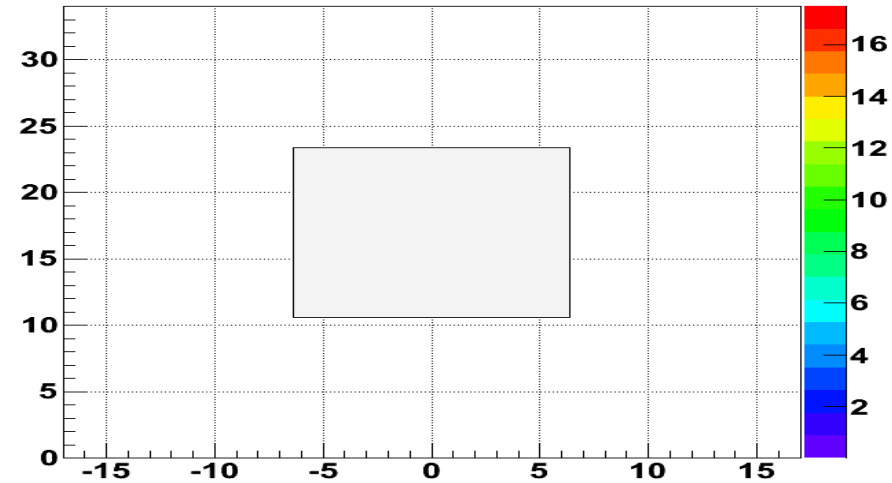
Simulated and Reconstructed Drell Yan Event

Event 1

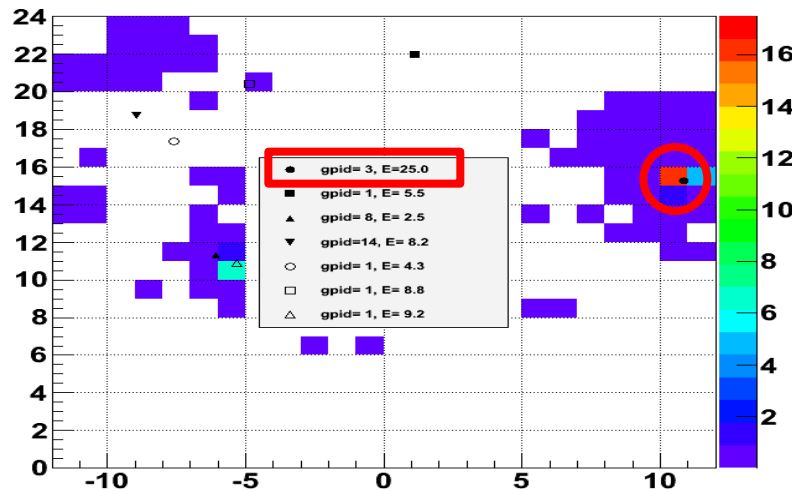
Geant energy



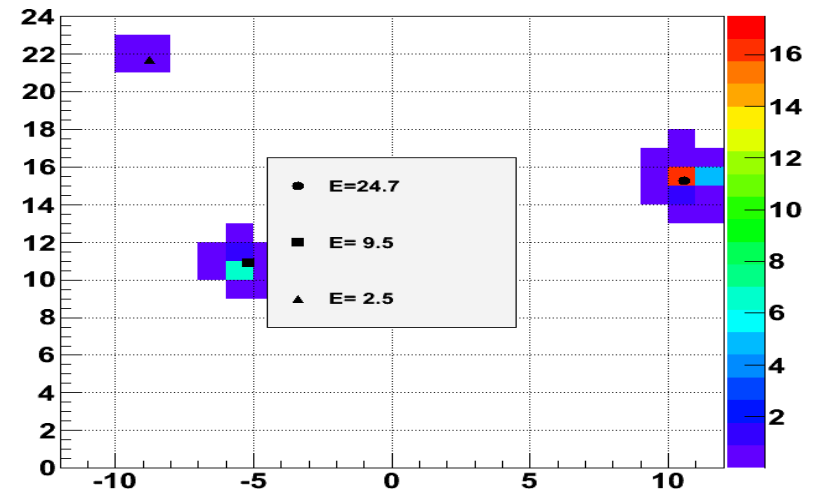
Fitted energy



Geant energy



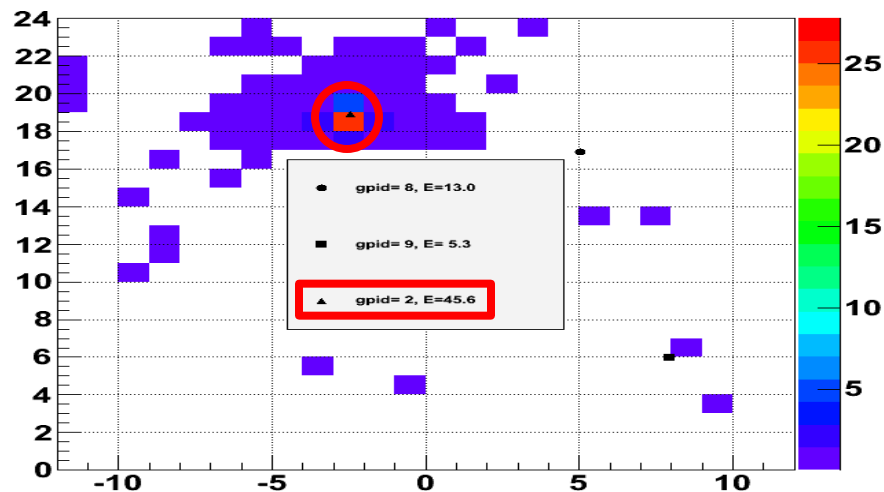
Fitted energy



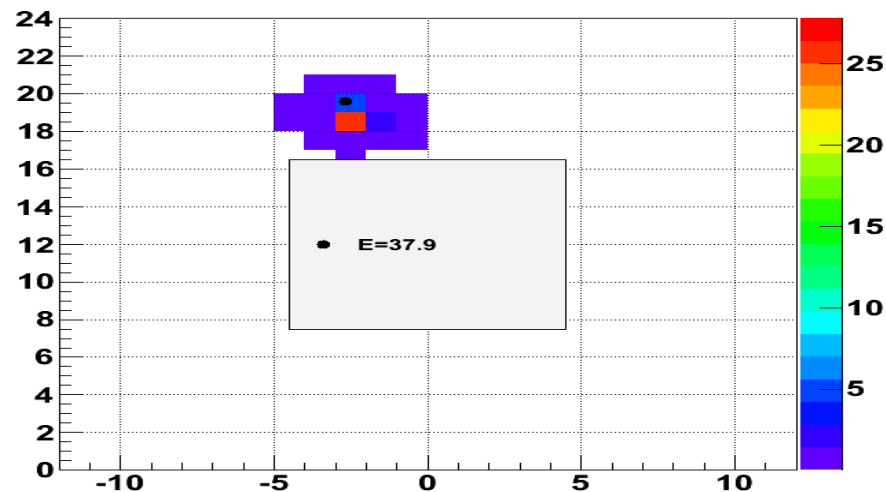
$$M = 3.457$$

Event 6

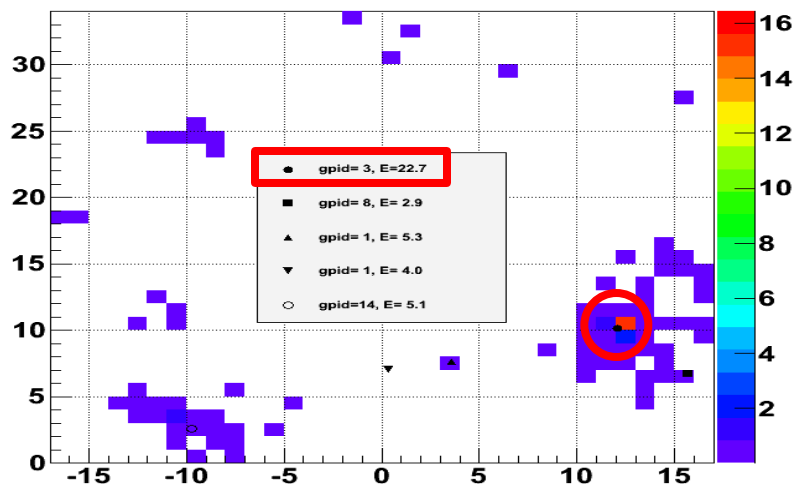
Geant energy



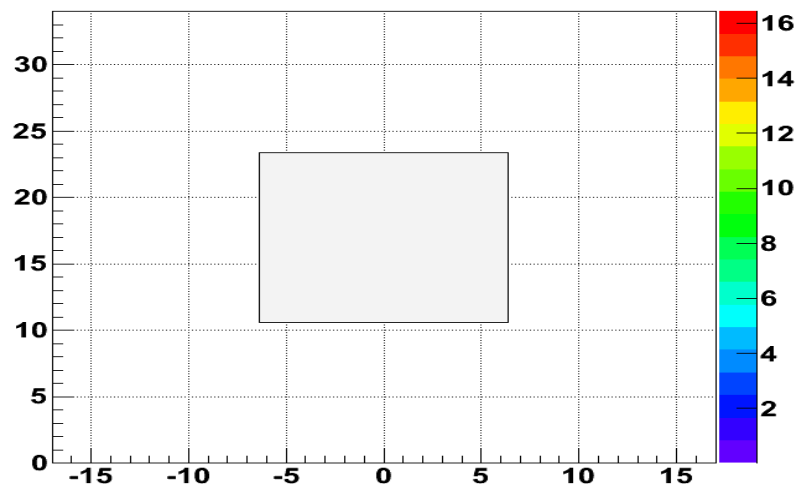
Fitted energy



Geant energy



Fitted energy

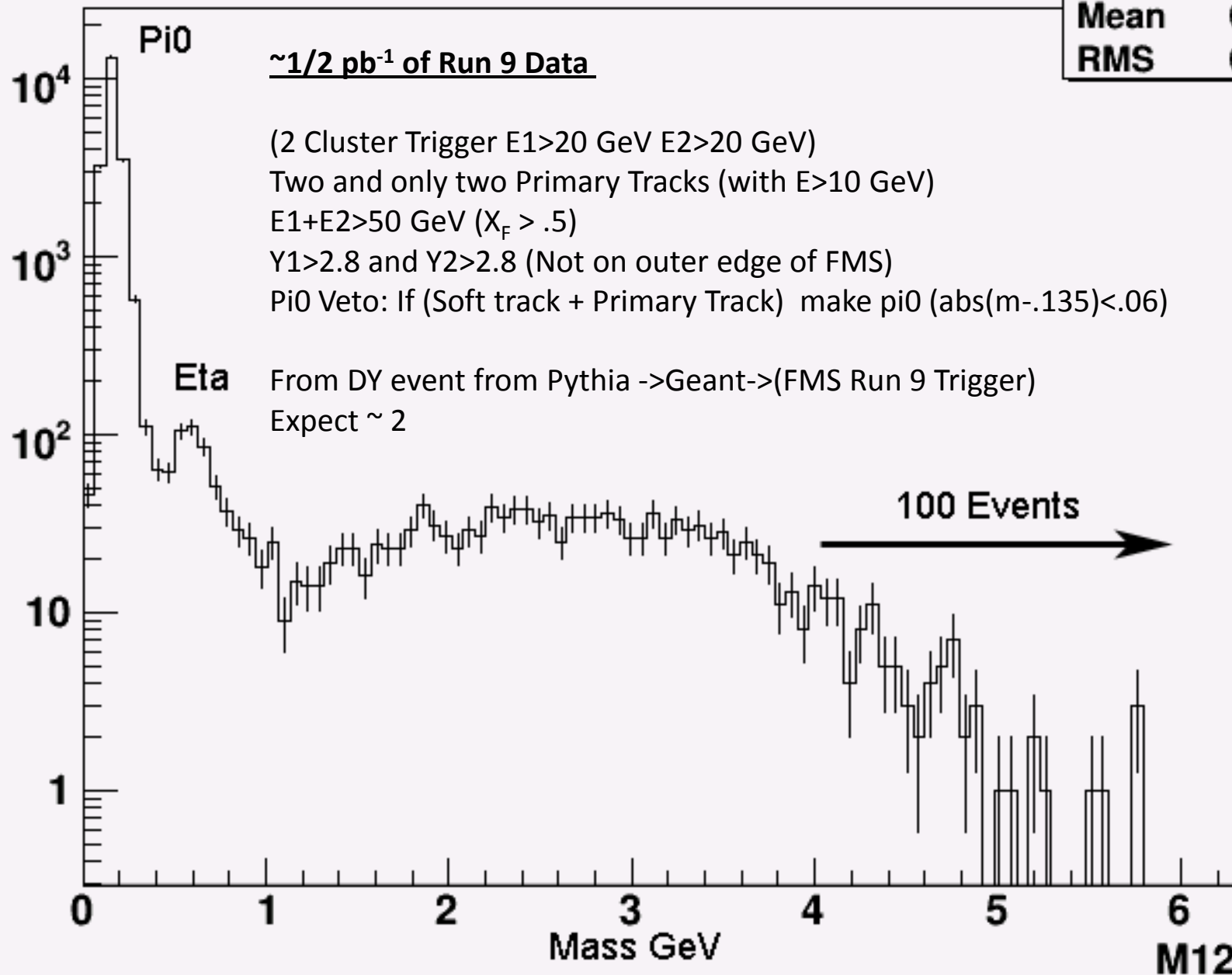


M = 4.128

M12 {(N12==2) &&E12>50 &&E12<80 &&E12*(1+Z)/2>20&&E12*(1-Z)/2>20&& abs(M12-6)<6 &&Y1>2.8 && Y2>2.8}

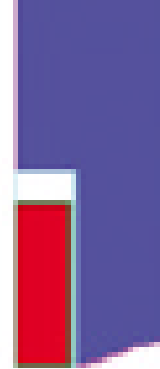
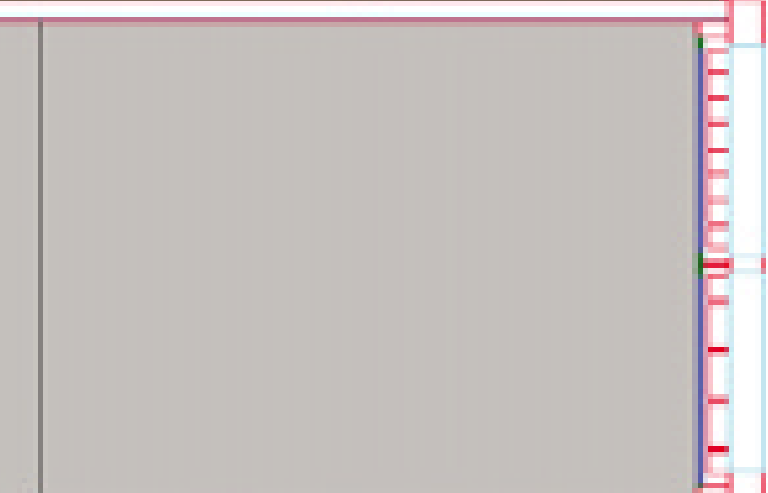
htemp

Entries	22564
Mean	0.3208
RMS	0.6318



Additional Cuts that could help for Run 9 Data for DY Candidates

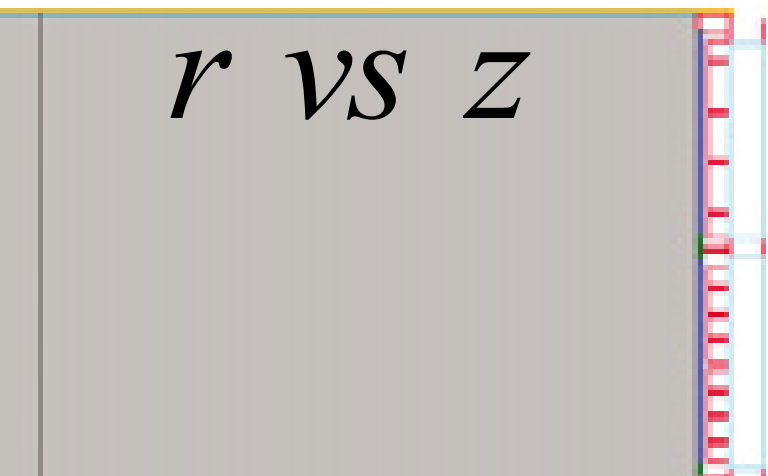
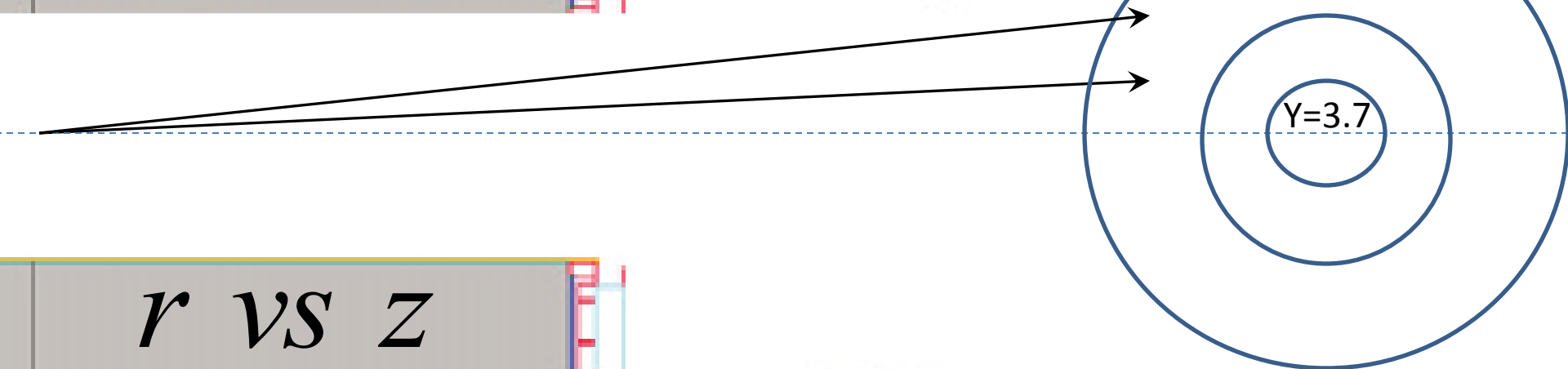
- Shower Shape to reject charged pions $(0.25 * 0.25) = 0.06$.
- Charge tracks to veto photon contribution
- Tracking inside magnet (see below) charged track signs for e^+ and e^- momentum match between momentum from magnetic curvature and EM.
 $(.1 * .1) = 0.01$ suppression of charged pions
- TRD electron id $(.1 * .1) = 0.01$ suppression of charged pions.



Magnet Hole

$Y=3.0$

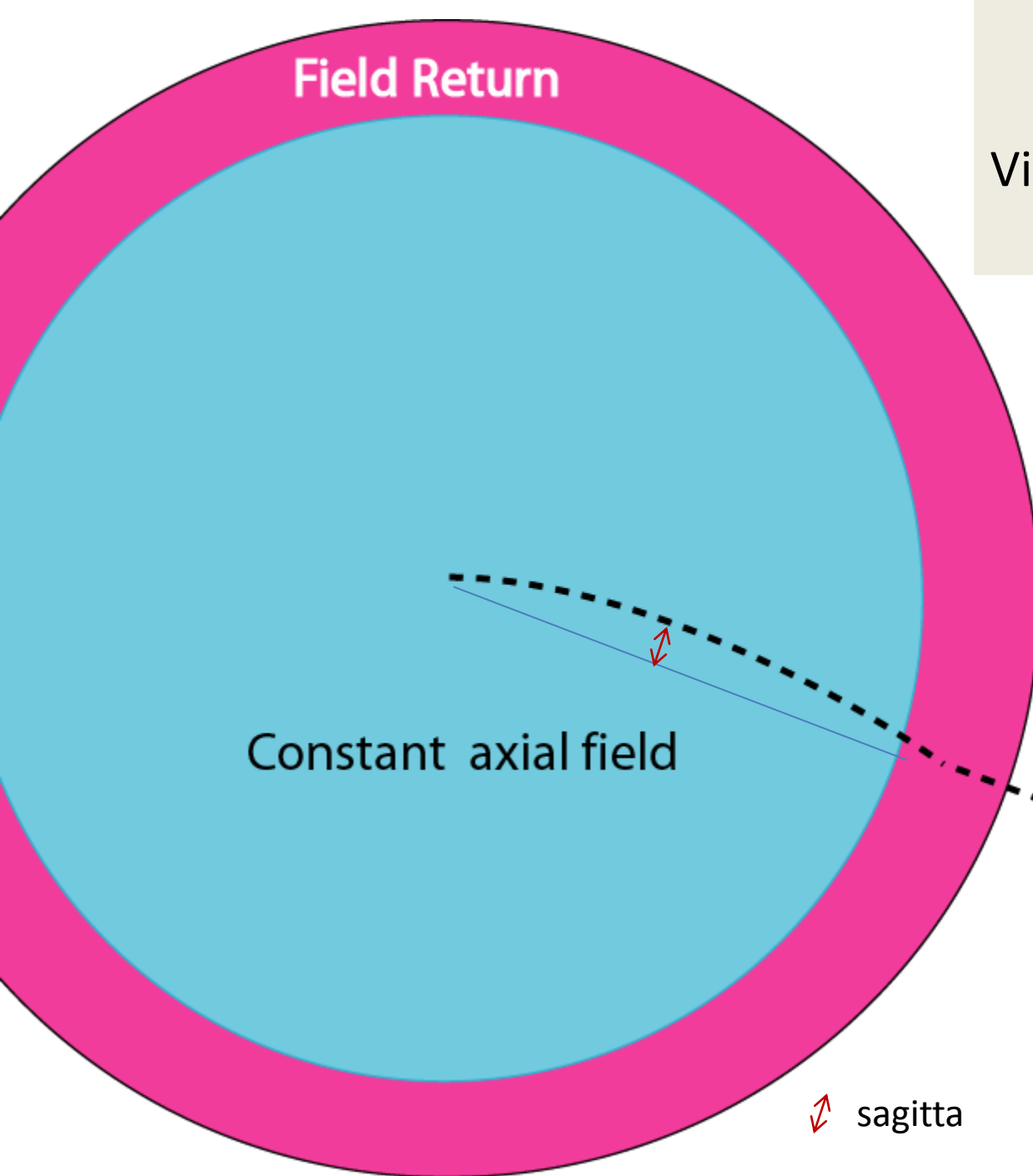
$Y=3.7$



x vs y

$$r \text{ vs } \phi$$

View of charged track in
magnetic field

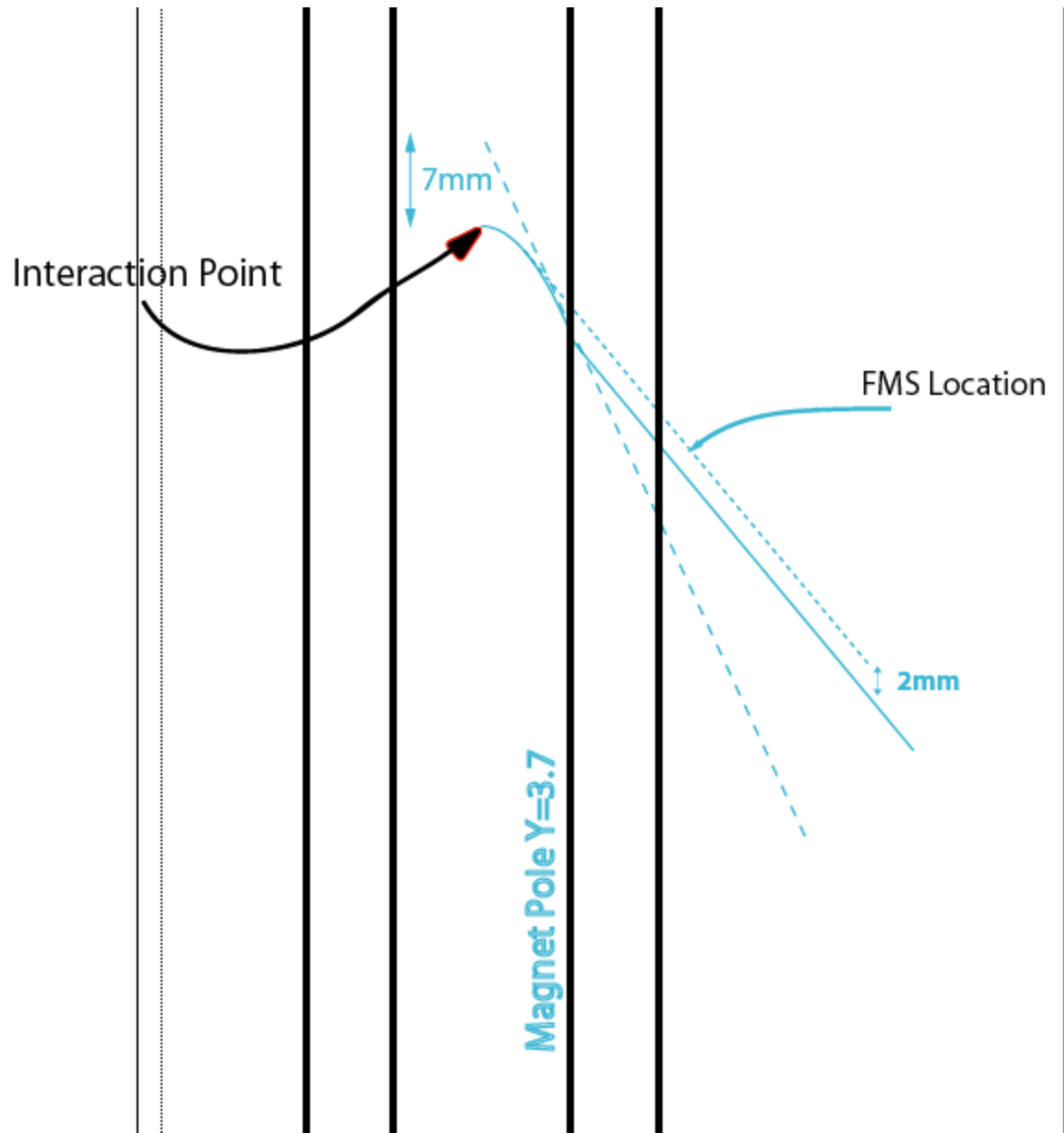


For a charged particle exiting
the constant field of a solenoid
magnetic field through the a
hole in a flux returning cap,

about $\frac{1}{2}$ of the angular
bend in the constant field
region is negated in the return
region.

sagitta

$$r \phi \text{ vs } r$$



For 20 GeV electron $Y=3.7$

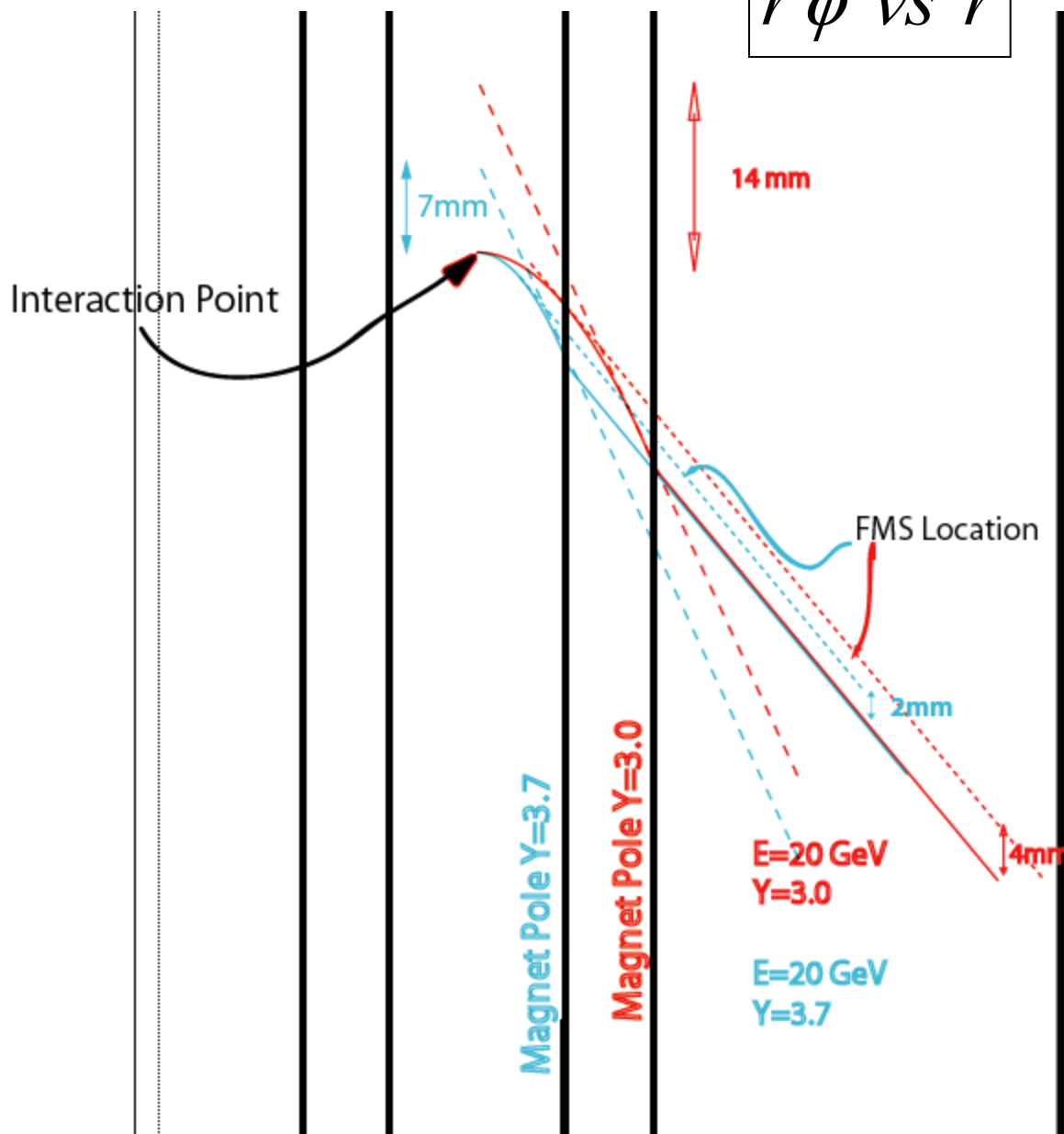
- Trajectory (r vs ϕ) measured just before exiting constant field

Projected back to vertex
Vertex displacement = 7 mm.

Projected forward to FMS ...
FMS displacement = 7 mm.

- Trajectory measured at origin
Projected to FMS ...
FMS displacement ~ 20 mm
- Trajectory measured at sagitta
Projected to FMS ...
displacement \sim sagitta ~ 1.5 mm

$$r \phi \text{ vs } r$$



For 20 GeV electron $Y=3.0$

- Trajectory (r vs ϕ) measured just before exiting constant field

Projected back to vertex

Vertex displacement = 14 mm.

Projected forward to FMS ...

FMS displacement = 14 mm.

- Trajectory measured at origin
Projected to FMS ...
FMS displacement ~ 40 mm

- Trajectory measured at sagitta
Projected to FMS ...
displacement \sim sagitta ~ 3 mm

Forward Magnetic Tracking Summary

A measurement of the charged track trajectory just inside the flux return (about 3 meters from the magnet center) can be projected to

- either the interaction vertex
- or to the FMS.

The displacement between the projected track and measured track will be called $D_\phi = r\Delta\phi$

$$\frac{\Delta p_t}{p_t} = \frac{\Delta p}{p} = \frac{\Delta D_\phi}{D_\phi}$$

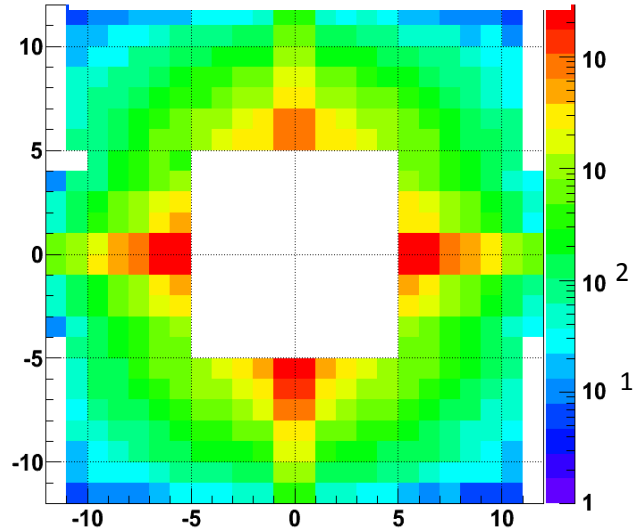
p_t, p	Y	For $\Delta D_\phi = 1mm$ $\frac{\Delta p}{p}$	For $\Delta D_\phi = 0.5mm$ $\frac{\Delta p}{p}$
1 GeV/c, 20GeV/c	3.7	15 %	8%
2 GeV/c, 40GeV/c	3.7	30%	15%
2 GeV/c, 20 GeV/c	3.0	7%	4%
4 GeV/c, 40GeV/c	3.0	15%	7%

Run 9 Setup Experience

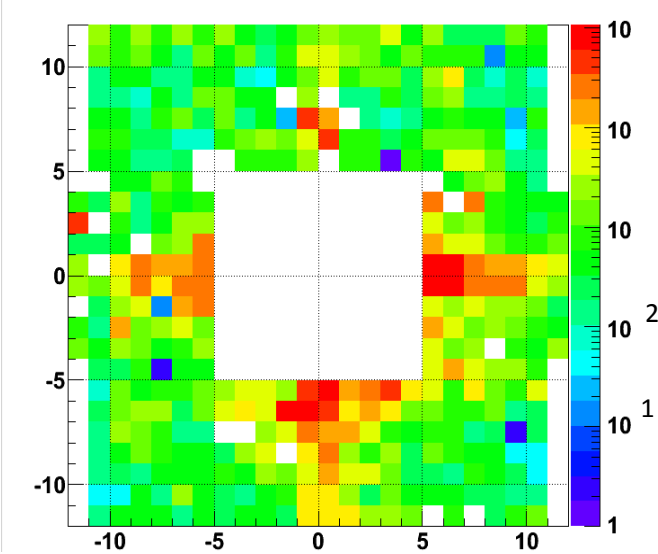
FMS Setup: Run 9 Experience (Jingguo Ma)

Simulated Trigger Rates
Ideal Gains (3 Triggers)

Trigger



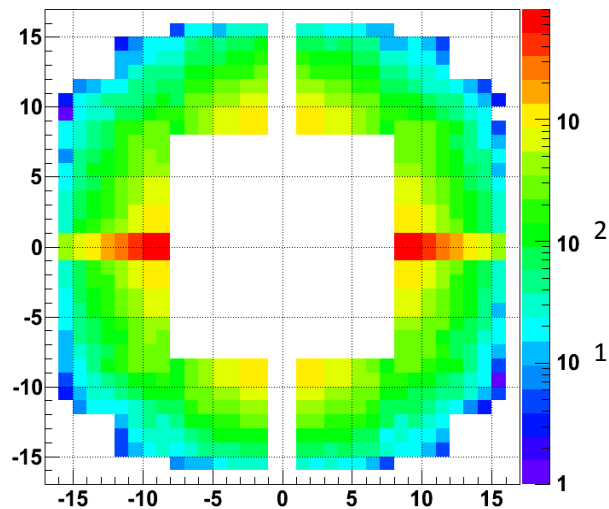
Run 9 Trigger Rates (analysis of Run 9 data)
Ideal Gains (3 Triggers)



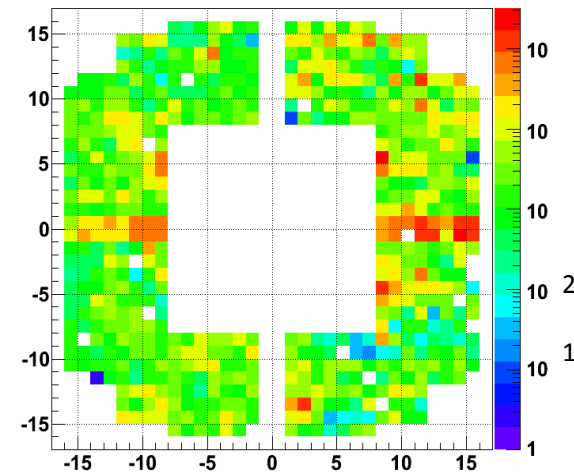
Small Cells

Fig. 7 Cluster sum trigger distributions for large cells with nominal gains. Most triggers are from the central horizontal region.

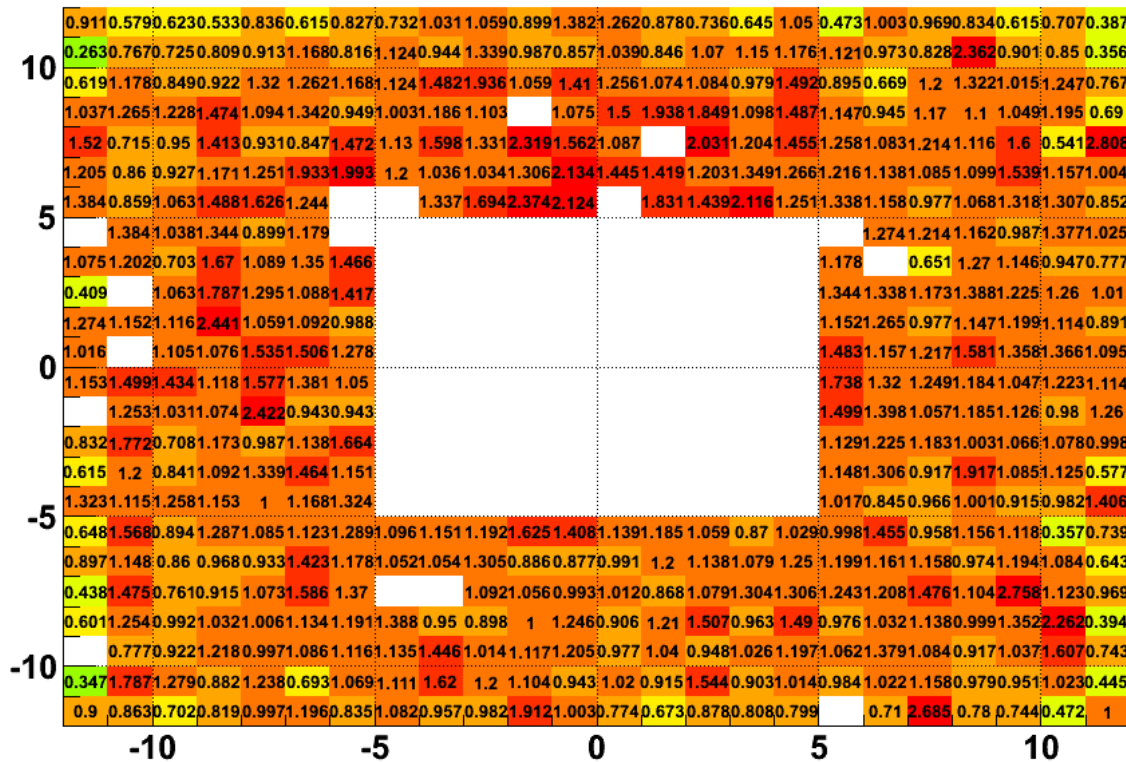
Trigger



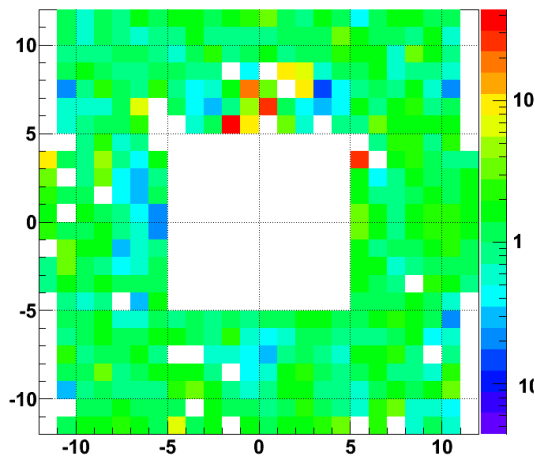
Trigger



Large Cells



Current Best FMS
Gain Correction

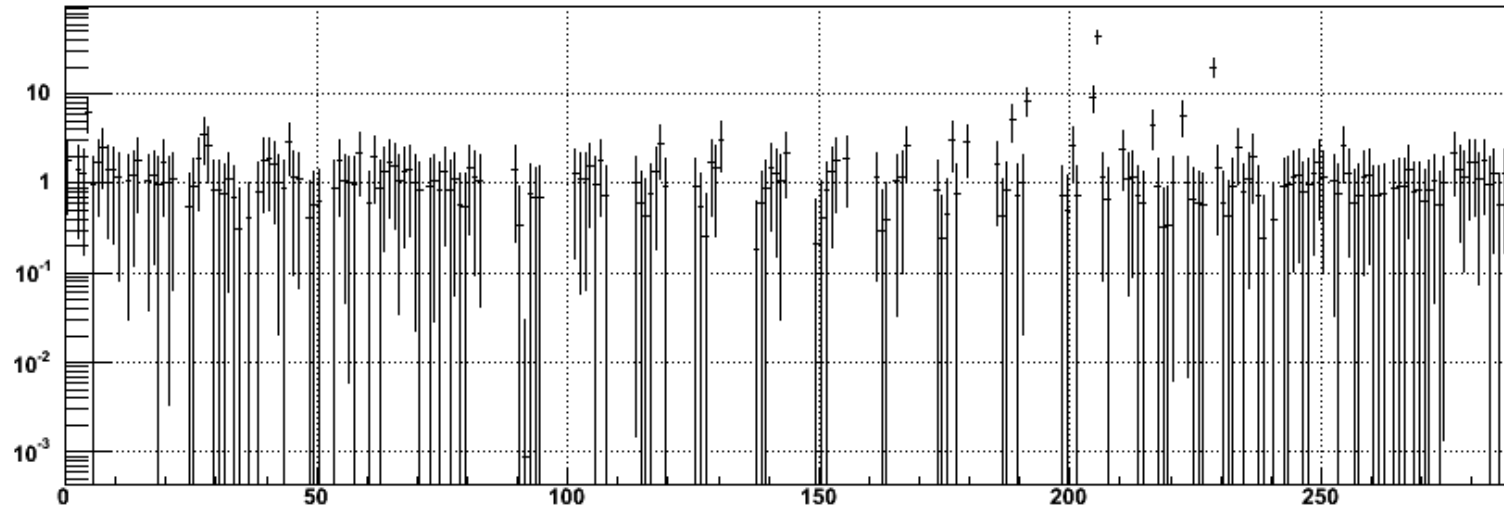


Ratio
Run 9 Trigger Rate to Simulated Rate

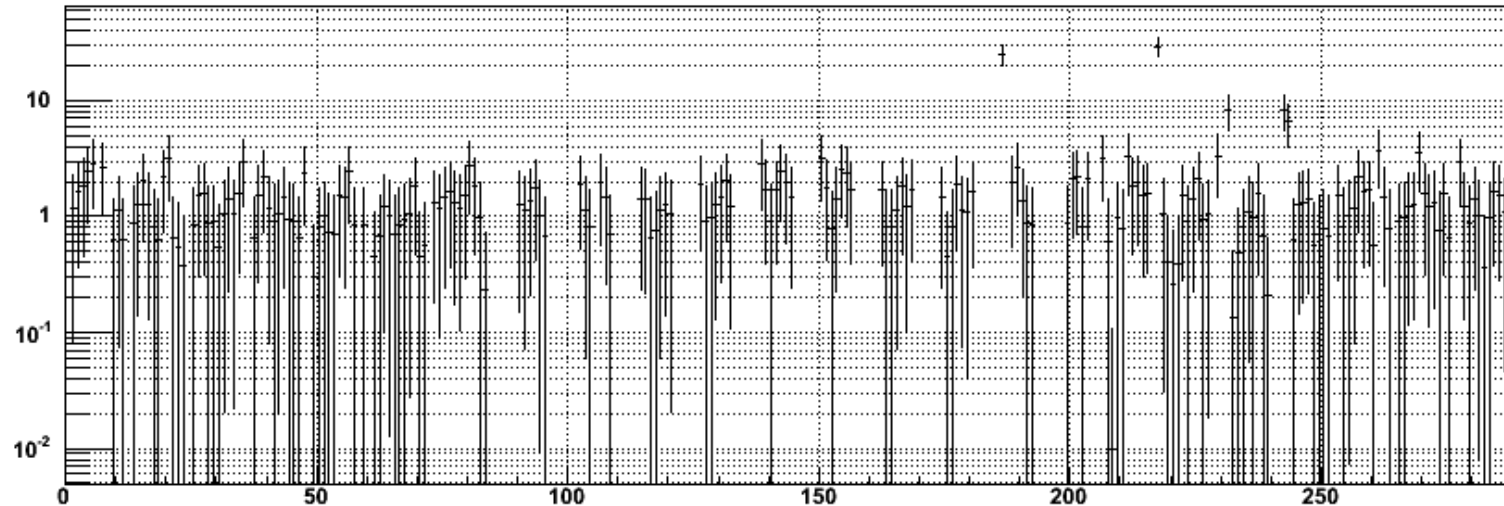
Ratio of Simulated Trigger Rates to Actual Rates

Sever Test of Calibration!

data to sim trigger rate ratio for north small



data to sim trigger rate ratio for south small



Goals:

- * well matched gains
- * FMS Gain Setup Without Pi^0 Reconstruction Iteration?

Consequence: Real Time HV
adjustment:

Fast: FMS Turn ON.